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MEASURED PLUME DISPERSION PARAMETERS OVER WATER VOLUME

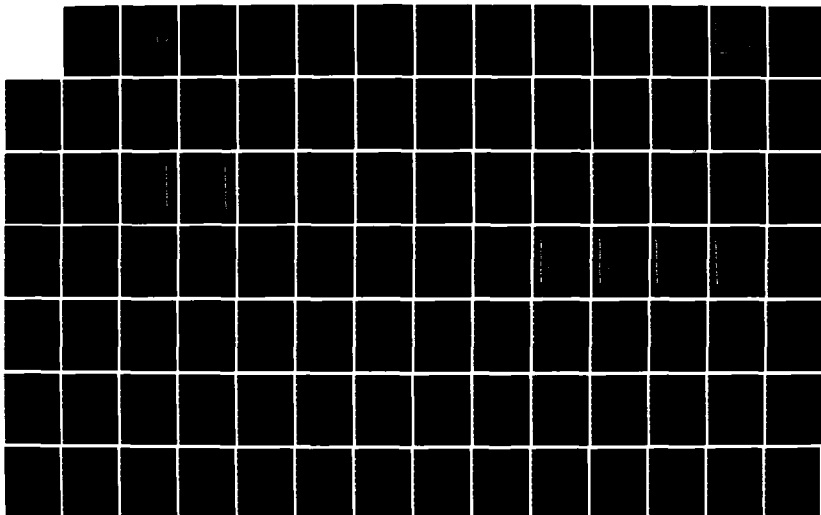
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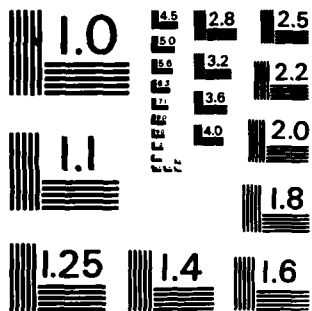
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MEASURED PLUME DISPERSION PARAMETERS OVER WATER  
VOLUME 1

by

C. E. Skupniewicz and G. E. Schacher  
SEPTEMBER 1984

Final Report for Period October 1983 - October 1984

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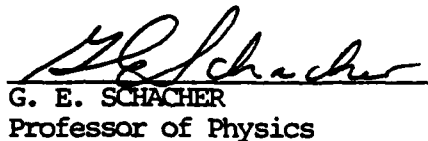
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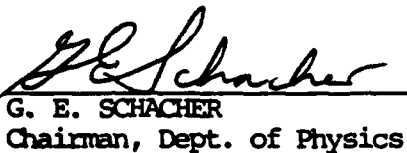
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Item 20. (continued)

is parameterized for the over water case by classifying the tracer results by stability in a Pasquill-Gifford equivalent scheme, and analytically describing horizontal and vertical plume growth as a function of plume travel distance. Several other over water data sets are used in this parameterization. Comparisons are made to the over land case.

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# MEASURED PLUME DISPERSION PARAMETERS OVER WATER

by

C. E. Skupniewicz and G. E. Schacher

## ABSTRACT

(cont R p 1)

*This volume analyzes the collected data.*

Data collected during a continuous, surface release, point source tracer experiment off the California coast is analyzed. The effects of high speed data collection from an airborne platform are removed by inverse transformation using the collecting instrument's transfer function. in frequency space. The tracer plume is characterized by a variety of parameters, including the conventional hourly averaged sigma-y and sigma-z values widely used in Gaussian plume dispersion formulae. Gaseous dispersion is parameterized for the over-water case by classifying the tracer results by stability in a Pasquill-Gifford equivalent scheme, and analytically describing horizontal and vertical plume growth as a function of plume travel distance. Several other over-water data sets are used in this parameterization. Comparisons are made to the over land case.

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## INTRODUCTION

→ The Minerals Management Service (formerly the Bureau of Land Management) sponsored a series of four atmospheric tracer experiments at California coastal locations over a two-year span, 1980-1982. These experiments were designed to assess air pollution impact from proposed oil exploration and drilling activities along the continental shelf. Two experiments (winter and summer) at each of two sites (open coast and Santa Barbara Channel) were funded in order to investigate air quality impact under a range of meteorological conditions and sites. A brief summary of these experiments and references is supplied in Table 1.

→ The basic designs of all four experiments were similar. A tracer gas, 100%  $\text{SF}_6$ , was continuously released from a stationary, sea surface platform located, for the majority of the experiments, approximately 3 miles from shore. During parts of the last experiment, the platform was moved to distances up to 5 miles from shore. <sup>(to p. 11)</sup> A variety of meteorological parameters were continuously monitored at various locations. Tracer gas concentrations were measured by a variety of methods at positions downwind of the release platform, with the majority made on or near the shore since the purpose was to assess potential on-shore air pollution impact. Experiments were limited to daytime periods of on-shore flow. Meteorological measurements, however, were not restricted to those time periods. This report utilizes a subset of the data base collected during the fourth experiments: offshore, aircraft,

Table 1

Central California Coastal Air

Quality Studies, 1981-1982

(Sponsored by Mineral Management Service)

DATE	LOCATION	FINAL REPORTS AVAILABLE	REF.*
Sep 80	Santa Barbara Channel Area	Aerovironment, Inc.	Zanetti et al. 81
Jan 81	"	"	"
Dec 81	Pismo Beach Area	Stanford Research Institute	Dabberdt et al. 83
Jun 82	"	Stanford Research Institute and Naval Postgraduate School	Dabberdt et al. 83 Other

\*Other reports available.

\*NPS work was sponsored by both the Minerals Management Service and the NPS Foundation Research Program.

continuous gas analyzer measurements.

The intention of this report is to characterize over-water diffusion from a continuous, near-surface, point-source release based upon these measurements. This report is built upon the meteorological results of Schacher et al. (1982) and a preliminary tracer gas and ranging results of Schacher et al. (1983). For a detailed description of the measurement techniques and data description, the reader is referenced toward these reports.

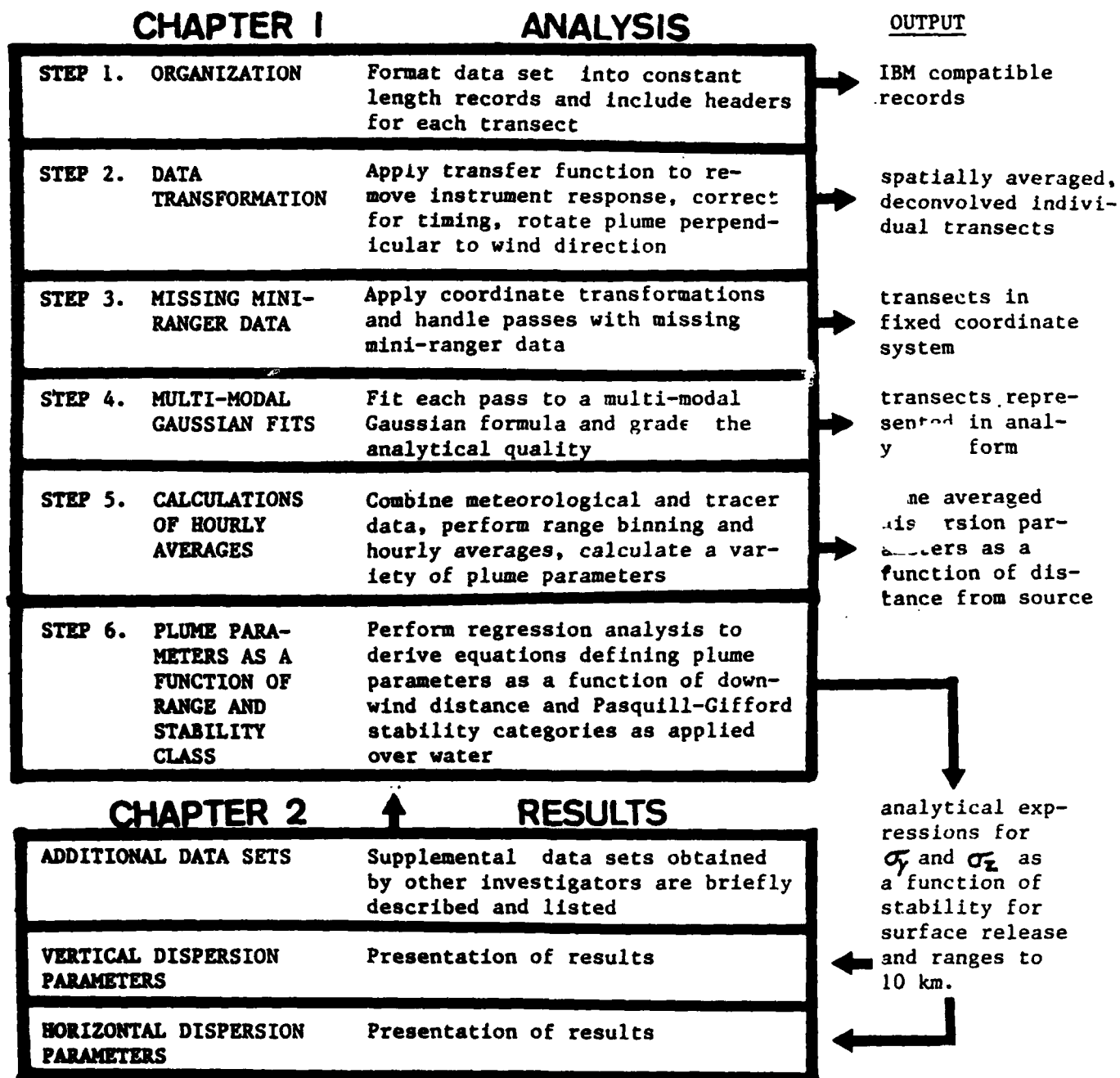
#### OUTLINE

A report flow chart is provided in Figure 1. This document is organized into two chapters with distinctly different designs. Chapter 1 contains technical procedures and data used in the piece-wise analysis of the data set. The second chapter presents one-hour average plume dispersion parameters,  $\sigma_y$  and  $\sigma_z$ , as a function of the well-known Pasquill-Gifford stability categories adapted for overwater use. Some additional data from other experiments supplement our data set to produce a more general parameterization.

Readers interested primarily in plume dispersion over water are advised to skip most of Chapter 1, concentrating mainly on Step 6 and Chapter 2. Those readers interested in the particular techniques used in the analysis of tracer data obtained from a high-speed platform may be more interested in Step 2. In addition to one-hour standard  $\sigma_y$  values, a wide variety of

Figure 1.

REPORT FLOW CHART



output is available from the Naval Postgraduate School Environmental Physics Group, and interested readers are advised to read Appendix A for a complete list of output data sets described throughout Chapter 1.

## CHAPTER 1 - DATA ANALYSIS

### Step 1 - Organization

The following data analysis was performed in a step-wise fashion, with the complete data set stored and cataloged in the Naval Postgraduate School's IBM 3033 mass storage system and 9-track tape at the completion of each step. Performing the overall analysis in six separate steps allowed for manual interrogation of the data set at each fundamental level and will allow for easy and flexible re-analysis of the data set in the future.

The analysis starts with  $\text{SF}_6$  concentrations, represented as digital voltage output from a continuous gas analyzer for single passes through the plume. Aircraft position was recorded from dual miniranger signals, resulting in paired position/ concentration data. Each plume transect was chosen to start where the analyzer first sensed  $\text{SF}_6$  along the flight path.

The original data set consists of seven files with one experimental day per file. Each file contains a different number of passes. Each pass starts a new (2048-byte) data block; the number of blocks needed depends on the pass length. Records are of variable length.

This data set was written into mass storage on the IBM 3033. The type of mass storage file used for this analysis is called a "partitioned data set". This data set consists of a number of user-specified "members". Each member can be accessed interac-

tively or via program control. If the members are less than 5000 lines, they can be edited interactively by the user. This was desirable; therefore, care was taken to keep each member under this limit. Also, members must consist of 80-character records. Therefore, the initial records became unsynchronized with the mass storage records after the transfer.

At this point, a simple program named FORMAT converted the variable length record format to a fixed length format. The output was interrogated and calibration passes<sup>†</sup> removed. Calibration factors derived were added to the header of each pass. These data were written to a partitioned data set named AIR2, residing on the Environmental Physics Group's private mass partitioned storage volume. All data set member names, format, and content will be presented in tabular form later in this report.

Next, the data set AIR2 was transformed to AIR3 by the program REDUCE. This program performed 3 vital functions. First, it converted raw voltages (corrected for background SF<sub>6</sub> concentration) to parts per trillion (PPT) concentration via the calibration factors mentioned above and experimentally-derived calibration formulae. The calibration factors were periodically measured during the experiment. The conversion formulae account for instrument non-linearity at high concentrations. The equations are:

---

<sup>†</sup> During a calibration pass, the instrument was purged with a "span" gas of known concentration in the instrument's linear region to obtain calibration factors.

$$V = \frac{V_B - V_0}{C} \quad (1)$$

where V is voltage normalized to laboratory conditions;

$V_0$  is output voltage from the analyzer;

$V_B$  is baseline (background) voltage;

C is the calibration factor determined during the experiment (See Table 2).

$$SF_6 = 5340V \quad [V < 1.345] \quad (2a)$$

$$SF_6 = \exp (1.160V^2 - 2.455 V + 10.122) \quad [1.345 < V < 1.687] \quad (2b)$$

$$SF_6 = \exp (1.461V + 6.823) \quad [1.687 < V < 2.053] \quad (2c)$$

$$SF_6 = \exp (4.252 V^2 - 16.780 V + 26.369) \quad [V > 2.053] \quad (2d)$$

$SF_6$  is concentration in parts per trillion.

Table 2  
Calibration Factors For Continuous SF6 Gas Analyzer

<u>Date</u>	<u>Time Period (PDT)</u>	<u>C</u>
6/21/82	BEGIN - 1640	.665
	1640 - END	.685
6/22/82	BEGIN - 1720	.695
	1720 - END	.685
6/24/82	BEGIN - END	.635
6/25/82	Begin - 1300	.620
	1300 - 1345	.615
	1345 - 1440	.605
	1440 - 1520	.600
	1520 - END	.615
6/27/82	BEGIN - 1720	.640
	1720 - END	.650
6/28/82	BEGIN - END	.670
6/29/82	BEGIN - 1620	.630
	1620 - END	.636

The second vital function performed by "REDUCE" was to determine plume transect Cartesian coordinates. This was accomplished, in most cases, with the mini-ranger data. Three scenarios existed, depending on mini-ranger performance for a given pass. When both mini-ranged distances were available, polynomial fits were performed to eliminate data "jitter" and simple triangulation used to determine plume coordinates. When one, or both, mini-ranger signals were intermittent, regression analysis was used where possible, to fill in the "gaps". When one or both mini-ranger signals were missing, coordinate determination was postponed for later analysis. An in-depth discussion of the above process design is given in Schacher et al. (1983).

The sampling grid coordinate system is shown in Figure 2. The mini-ranger transmitters were located on the beach, and north and south buoys were located so as to aid aircraft navigation. Their grid map locations, along with the variable ship locations are given in Table 3.

Figure 2.

# NPS SF<sub>6</sub> TRACER STUDY SAMPLING GRID

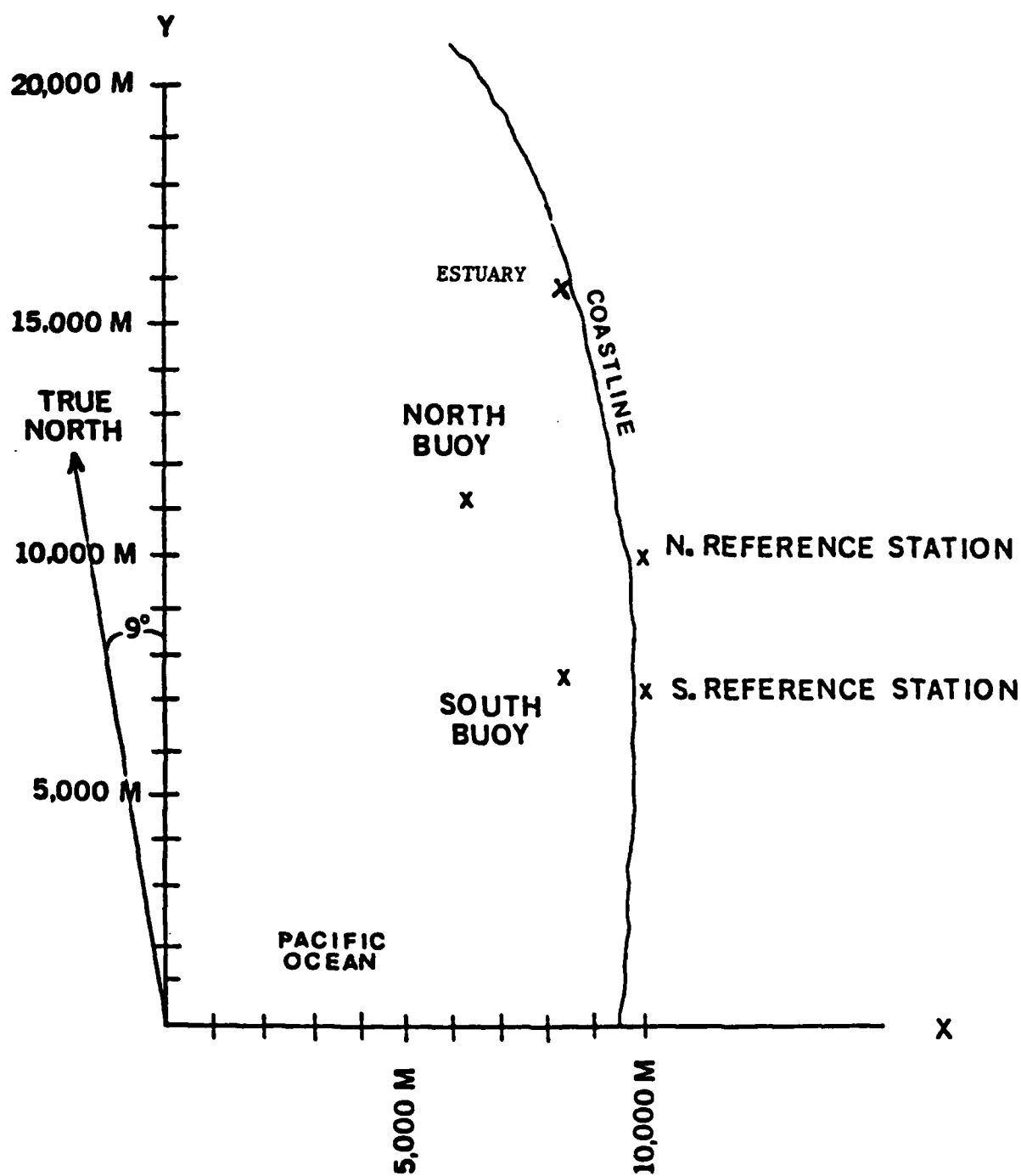


Table 3  
Grid Map Locations

Reference (See Figure 2)	Grid Coordinates (meters)	
	X	Y
N. Bouy	5926	11140
S. Bouy	8114	7550
Ship 6/21/82	4055	10200
" 6/22/82	4945	6369
" 6/24/82	4103	8628
" 6/25/82	4111	8601
" 6/27/82	399	11090
" 6/28/82	581	12493
" 6/29/82	1120	10459
Estuary	8896	15430
N. Ref. Station	10000	10000
S. Ref. Station	10000	7050

Figure 3a.

MASS STORAGE DATA SET "AIR2";  
ABBREVIATED SAMPLE OF ANALYSIS STEP 1 OUTPUT

FILE: AIR2.DAT

```

P 7/24-14.35:17 1 1000
B 7.519 474 277 1000
D 0 7155 -0.017
D 1 4757 -0.015
D 0 7155 -0.015
D 1 5757 -0.015
D 0 7157 -0.015
D 1 4777 -0.015
D 0 7157 -0.015
D 1 4771 -0.015
D 0 7157 -0.015
D 1 4775 -0.015
D 0 7159 -0.011
D 1 5775 -0.017
D 0 7159 -0.015
D 1 5777 -0.015
D 0 7159 -0.015
D 1 4777 -0.015
D 0 7159 -0.015
D 1 4777 -0.015
D 0 7159 -0.015
D 1 5777 -0.015
D 0 7159 -0.015
D 1 5759 -0.015
D 0 7159 -0.015

```

•  
•

INDIVIDUAL LINE KEY: column one is the code

P: date, time, pass number, data quality index, altitude  
B: total elapsed time (sec), total number of points,  
mini-ranger 0 number of points, mini-ranger 1 number  
of points, number of bad points, calibration factor  
D: mini-ranger code, distance from mini-ranger (m),  
analyzer output (volts)

Figure 3b.

MASS STORAGE DATA SET "AIR3";

ABBREVIATED SAMPLE OF ANALYSIS STEP 1 (continued)

```

1  7211410131 13.51 1.000 1.000
2  1.45655 1.0141211 1.000 1.0141211 1.000
3  1.11111 1.0141211 1.000 1.0141211 1.000
4  1.11111 1.0141211 1.000 1.0141211 1.000
5  1.11111 1.0141211 1.000 1.0141211 1.000
6  1.11111 1.0141211 1.000 1.0141211 1.000
7  1.11111 1.0141211 1.000 1.0141211 1.000
8  1.11111 1.0141211 1.000 1.0141211 1.000
9  1.11111 1.0141211 1.000 1.0141211 1.000
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100 1.11111 1.0141211 1.000 1.0141211 1.000

```

INDIVIDUAL LINE KEY:

- lines 1-7: header information (self explanatory)
- 8-11: mini-ranger statistics
- 12: R/V ACANIA position in grid (see Figure 2)
- 14-: data -time in seconds
  - width is distance from first non-zero concentration in meters
  - concentration in PPT

## Step 2 - Data Transformation

This step accounts for instrument effects on the data. If the data are perceived as a time series, and we treat the instrument as a first-order linear system\*, then,

$$\frac{d\tilde{X}}{dt} = \tilde{A}\tilde{X} + \tilde{B}\tilde{U} \quad (3)$$

where  $\tilde{X}$  is a one-dimensional matrix of state variables;

$\tilde{U}$  is a matrix inputs;

$\tilde{A}, \tilde{B}$  are square matrices of coefficients; and

$t$  is the independent variable.

In general, the system output can be represented as a linear combination of the state variables and the inputs. In this specific case, the input is the true  $SF_6$  concentration, and the output of interest is a state variable; the measured  $SF_6$  concentration. Also, this case is concerned with only one state variable; therefore, matrix expressions will be dropped. A solution can be expressed as the convolution of the input waveform and a function called the unit impulse response of the system.

$$x(t) = \int_{-\infty}^{+\infty} h(t - \tau)u(\tau)d\tau \quad (4)$$

where  $x(t)$ ,  $u(t)$  are singular state and input variables,  
respectively;

$h(t)$  is the unit impulse response.

\*In a second-order system, a second state variable would simply be the derivative of the first variable.

Convolutions are rather difficult to perform on digital machines; therefore, we use the convolution theorem, which states that convolution in the time domain is analagous to multiplication in the frequency domain.

$$x(t) = F^{-1}[X(f)] = F^{-1}[H(f)U(f)] \quad (5)$$

where  $X(f)$  is the Fourier transform of  $x(t)$

$U(f)$  is the Fourier transform of  $u(t)$

$H(f)$  is the Fourier transform of  $h(t)$ , or the "transfer function"

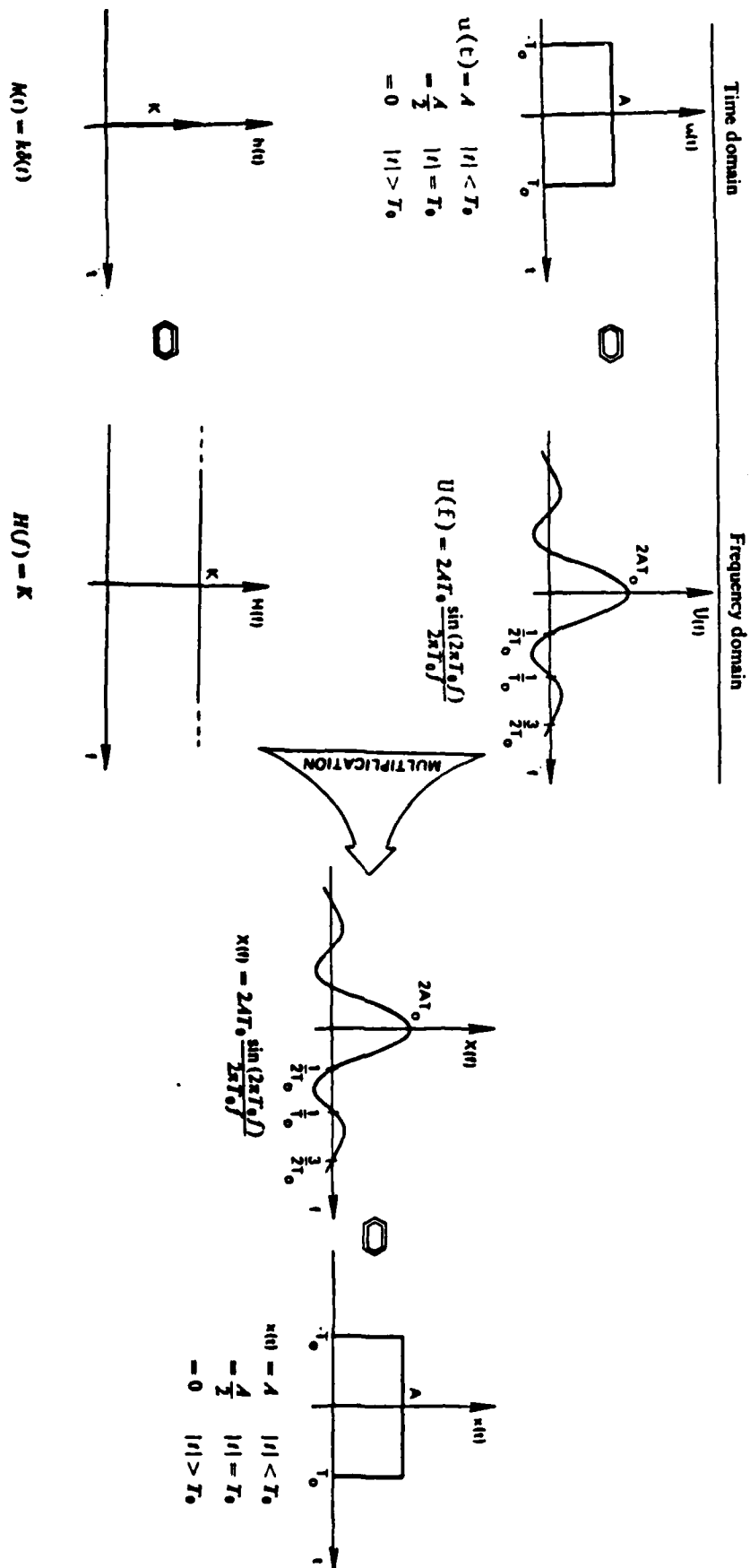
$F^{-1}$  refers to the inverse Fourier transform.

Finally, since the system input is the desired quantity, Equation (5) is inverted, yielding:

$$u(t) = F^{-1}[X(f)/H(f)] \quad (6)$$

A graphical example can serve as a "proof" of this concept, often referred to as the transfer function approach (see Figure 4). Let the unit impulse response be the unit impulse. The impulse transforms to a constant function of magnitude 1, while  $x(t)$  transforms to  $X(f)$ . Their product is identically  $X(f)$ , and the inverse Fourier transform yields  $x(t)$ . It is obvious that any input function,  $u(t)$ , will produce an output,  $x(t)$ , identical to itself, as it should, if the system is transparent. This example should not be considered complete proof of the transfer function approach, but merely demonstrates an extreme situation.

Figure 4. GRAPHICAL EXAMPLE OF TRANSFER FUNCTION APPROACH AND "PROOF"



[If  $h(t)$  is the unit impulse,  $K=1$ ]

The program developed to apply the transfer function was called XFORM. The Cooley-Tukey fast Fourier transform routine was used as the core of this program. No tapers were applied to the time domain truncation function in order to reduce leakage because the frequency distribution of the waveform was unknown. All high frequency information was desirable, and a taper could have destroyed that information. Also, Hanning or cosine windows often smooth the waveform. This would artificially widen the plume; an undesirable effect. To keep computations to a minimum, the number of points in the discrete Fourier transform should be small. Crow and Tewscher (1983) determined the proper number to be 18, based on the instrument high frequency cut-off and the approximate airspeed. The program XFORM therefore averaged each pass into an 18 point series before applying the transform. Each point, therefore, represented upwards of 100 samples. If the measurement variability between samples is considered independent, this would decrease the statistical significance of measurement errors tenfold. Considering the nature of the noise (instrument noise, intake airflow dynamics, etc.) and the errors produced, the data density achieved in this experiment appears to be more than necessary to achieve sufficiently small error. Ten to twenty samples per data point would have produced accuracy to within 50 ppt, an acceptable level. The 18 point series was designed so that the records start and end at zero concentration, with all other points non-zero, to avoid introducing false high-frequency components due to discontinuity or background noise. The

untransformed data set was stored on mass storage for comparison to the transformed data

The program XFORM next entered the transfer function subroutine. The first task in the subroutine was to determine the transfer function. This was accomplished by first transforming the experimental time series; a simulated "unit impulse" as the input waveform, and the resultant measured SF<sub>6</sub> concentration as the output. Next, each frequency component's contribution to the transfer function was determined by dividing the input by the output. As implied in the earlier discussion, using an impulse as input produces a smooth function in frequency space, contributing information to the transfer function from all frequency components. Since the results of the Fourier transforms are imaginary numbers, their quotient is also imaginary, as follows:

$$H(f)^{-1} = \frac{Y(f)}{X(f)} = \frac{(a_1 + b_1 i)}{(a_2 + b_2 i)} = \frac{a_1 a_2 + b_1 b_2}{a_2^2 + b_2^2} + \frac{(b_1 a_2 - a_1 b_2)}{a_2^2 + b_2^2} i \quad (7)$$

where  $H(f)^{-1}$  is the inverse transfer function;

$Y(f)$  is the transform of the laboratory "impulse";

$Z(f)$  is the transform of the laboratory output;

$a_1, a_2$  are the real parts of the input and output transforms;

$b_1, b_2$  are the imaginary parts of the input and output transforms.

The experimental time series (measured output) is then transformed, and multiplied by the inverse transfer function to yield the input waveform,

$$U(f) = X(f)H(f)^{-1} = \frac{a_1 a_2 a_3 + b_1 b_2 a_3}{a_2^2 + b_2^2} - \frac{b_1 a_2 b_3 - a_1 b_2 b_3}{a_2^2 + b_2^2} + \left[ \frac{a_1 a_2 b_3 + b_1 b_2 b_3}{a_2^2 + b_2^2} + \frac{b_1 a_2 a_3 - a_1 b_2 a_3}{a_2^2 + b_2^2} \right] i \quad (8)$$

where  $U(f)$  is the transform of "true" input waveform;

$X(f)$  is the transform of measured output waveform;

$a_3$  is the real part of the output transform;

$b_3$  is the imaginary part of the output transform.

Finally, an inverse transform yields the "true" input time series.

XFORM next called the DELAY and ROTATE subroutines. These subroutines operated on the coordinates of the pass; therefore, when no navigation information was available, they were not used. The DELAY subroutine applied a constant time delay, translated as a shift in the coordinates, to account for the lag time created by system dynamics. This lag time was obtained daily in situ tests. The ROTATE subroutine corrected the concentrations to produce values appropriate to a pass perpendicular to the mean wind direction. In almost all cases, this correction proved to be negligible, since the flight paths were usually within 5 degrees of the desired direction.

The two resultant mass storage data sets were called AIR4 (untransformed passes) and AIR5 (transformed passes). Figures 5 and 6 show two examples of untransformed and transformed data. The abscissa represents distance from first detection of  $\text{SF}_6$ . The apparent change in the peak location upon transformation results from the inherent time shift due to the "smoothing" of the input waveform. In Figure 5, the plume has been significantly narrowed, and the mass conserved with an increase in peak concentration. Also, a second mode appears which corresponds to a slight inflection in the untransformed data. This demonstrates the usefulness of the transfer function approach for retrieving high-frequency information. Figure 6 displays a much broader plume than in Figure 5. The transformation does not significantly change this waveform shape. Evidently, the transformed plume contains significant terms only at frequencies below those affected by the transfer function. Also note that the peak shift remains, since time shifting translates to phase shifting in frequency space, affecting all frequencies.

date : 6/28 time : 1415:49 pass 27

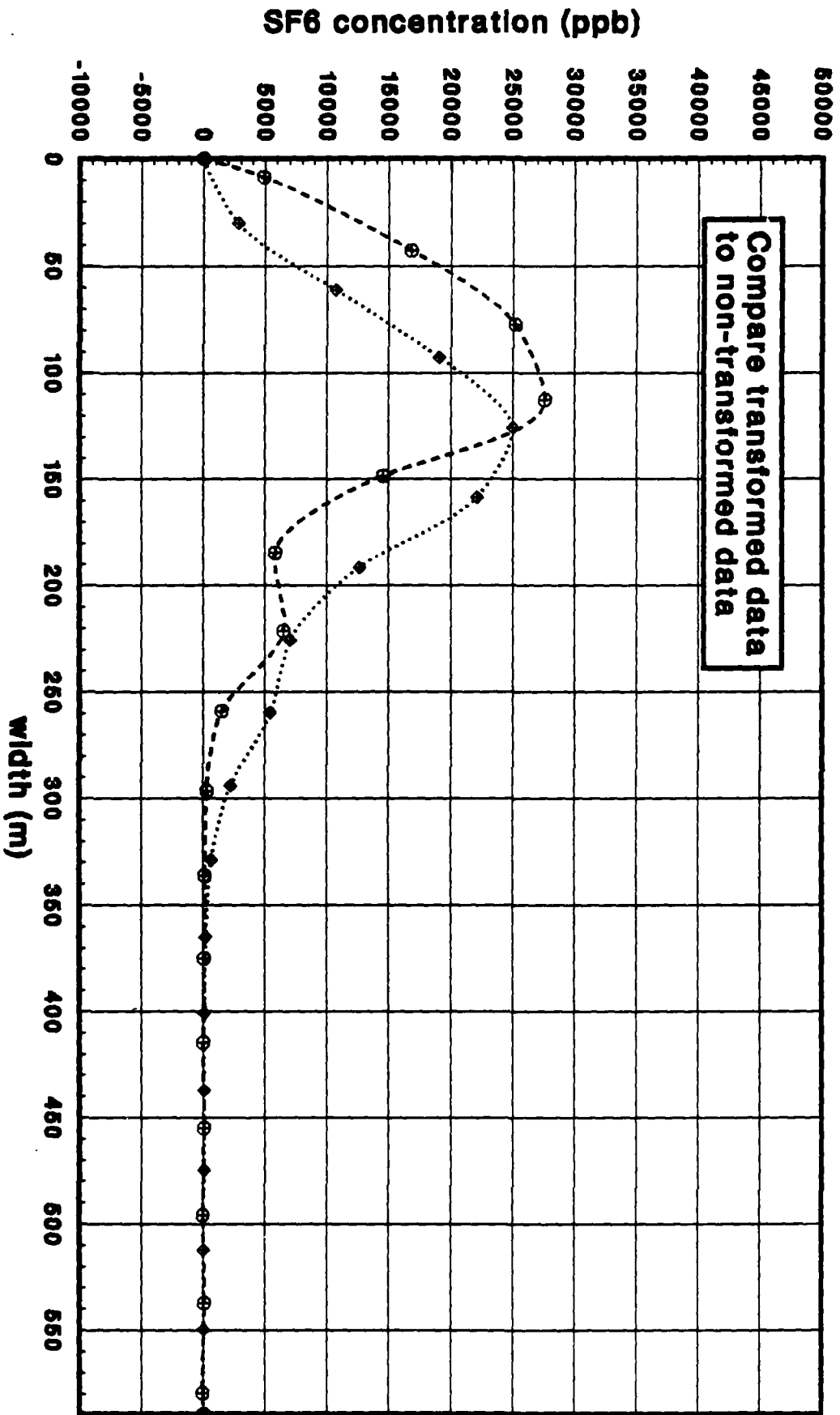


FIGURE 5. EXAMPLE OF ANALYSIS STEP 2 OUTPUT (NARROW PLUME)

**date : 6/28 time : 15 3: 4 pass 43**

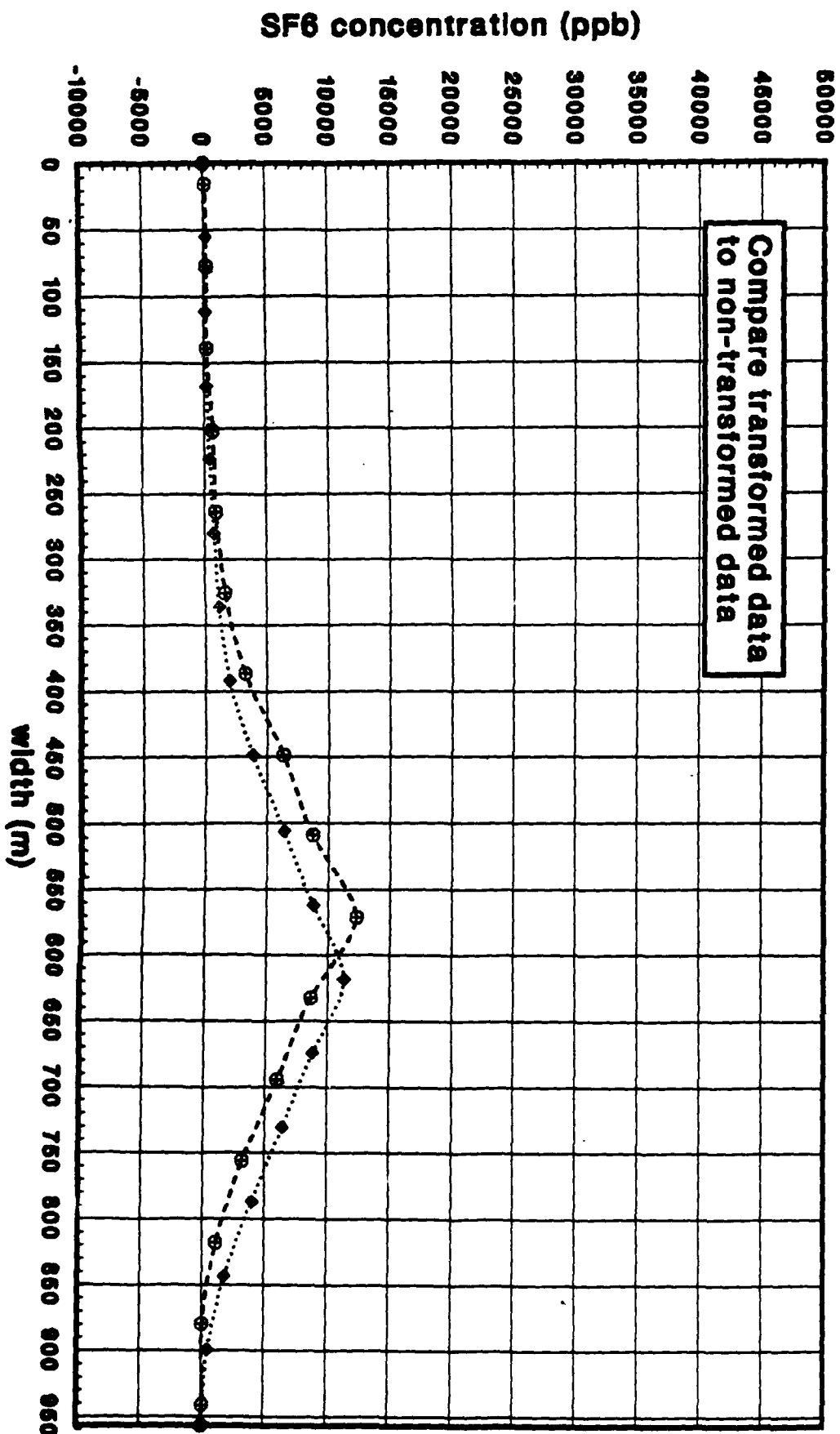


FIGURE 6. EXAMPLE OF ANALYSIS STEP 2 OUTPUT (WIDE PLUME)

### Step 3 - Missing Mini-Ranger Data

As mentioned in step one, at times during the experiment one or both mini-ranger signals were not available. Anticipating this problem, various reference points were selected as starting points for plume transects, and the time of intersection logged on the SF6 analyzer strip chart by an onboard technician. With this information, the flight heading, and airspeed, the plume coordinates could be estimated.

The program MINIFIX was written to do the necessary analysis. The heart of this program is a look-up table that lists passes with missing coordinates, and their associated reference point to plume peak distances. These distances were manually extracted from the SF<sub>6</sub> analyzer strip charts. Another table identifies the reference points for the various passes. These reference points are shown in Figure 2 and their positions listed in Table 3.

Since the reference point passage was logged instantaneously and the strip chart data was output by the analyzer, the inherent lag due to the system dynamics (mentioned in the previous section) had to be removed by MINIFIX. Care was taken to operate on the untransformed plume when applying this correction, to avoid the time-shift due to the transformation process.

As in the previous section, the plume was rotated to a wind direction perpendicular alignment. The untransformed and transformed data sets were output to mass storage files AIR6 and AIR7 respectively. Table 4 lists all SF6 cross-sections archived at

NPS, and also identifies those passes with missing mini-ranger data. An example of the data sets, AIR4 - AIR7 (identical formats) is presented in Figure 7.

Table 4

Pass Numbers for Which Plume Cross Sections Were Determined\*

(complete analysis through the "Missing Mini-Ranger" Step)

DATE (June 82)

21	22	24	25	27	28	29
54m	4m 58 m	1 41	1 36	1 42	1 67	1 56
55m	5m 59 m	2 42	2 37	2 43	3 68	2 57
56m	6m 62 m	3m 43	3 38	44	4	3 58
8m 57m	7m 63 m	45	4 39	4 46	6 71	4 59
9m 59m	8m 64 m	7 46	6 41	5 47	9 72	5
	62m 11m 65 m	8 47	7 42	7 48	11 73	7
11m 63m	12m 34 m	9 48	8 43	8 49	13 74	12
13m 64m	13m 42 m	10 49	9 44	9 50	15 77	14 19m
14m 65m	14m 3 m	12 53	10 45	51	78	18 22m
15m	15m 26m	13 54	11 46	11 52	17	23 24m
17m 68m	16m 27m	14 55	12 48	14 53	19	27
19m 72m	17m 30m	15 56	13 49	16	22	29 26m
21m	18m 32m	16 57	14 50	17	26	31t
23m 75m	19m 61m	17 59	15 51	18	27 18m	32
27m 16m	21m 35m	18 61	16 52	19	21m	33
28m 18m	22m 36m	19 62	17 53	22	29 23m	
30m 25m	23m 37m	21 63	19 54	24	30 53m	55
32m 29m	25m	22 64	21 55	25	33 55m	37
34m 31m	29m	23 66	22 58	26	34	
36m 33m	31m	24 70	23 59	27	35 77m	39
37m 35m	33m	25 72	24 61	28	37 36	41m
38m 39m	41m	26 73	25	29		42
41m 42m	45m	28 74	26 63	30	44	43m
43m 58m	49m	29 75	27 64	31	46	44
	61m	30 76	28 65	32	48	45
45m 69m	51m	31 77	29 66	33	54	46m
46m 71m	52m	32 80	30 67	34	59	47
47m 73m	53m	33 81	31 68	35	60	48
48m 5m	54m	36	61	50		
49m	55m	37	33 70	37	62	51
	44m 56m	38	34 73	39	65	52
53m	57m	39	35 57m	41	66	55

pass numbers

- \* stored at NPS computer center as "AIR6" (untransformed), and "AIR7" (transformed)
- m mini-ranger not operating during this pass
- t this pass represented as two logical records

Figure 7.

EXAMPLE OF ANALYSIS STEPS 2,3 OUTPUTS

FILE: 1111      DAYS: 1      NAVAL POSTGRADUATE SCHOOL

DATE: 12/11/75      TIME: 14:11:17      PASS: 44  
 WIND: 10.0      DWD: 2.0      SF6: 10.0      SF6: 10.0  
 T-AVE: 10.0      T-STD: 1.0      T-STD: 1.0      T-STD: 1.0  
 S-F6: 10.0      S-F6: 10.0      S-F6: 10.0      S-F6: 10.0

TIME	10.0	10.0	10.0	10.0
0.0000	0.0	0.0	0.0	0.0
0.1250	0.0	0.0	0.0	0.0
0.2500	0.0	0.0	0.0	0.0
0.3750	0.0	0.0	0.0	0.0
0.5000	0.0	0.0	0.0	0.0
0.6250	0.0	0.0	0.0	0.0
0.7500	0.0	0.0	0.0	0.0
0.8750	0.0	0.0	0.0	0.0
1.0000	0.0	0.0	0.0	0.0
1.1250	0.0	0.0	0.0	0.0
1.2500	0.0	0.0	0.0	0.0
1.3750	0.0	0.0	0.0	0.0
1.5000	0.0	0.0	0.0	0.0
1.6250	0.0	0.0	0.0	0.0
1.7500	0.0	0.0	0.0	0.0
1.8750	0.0	0.0	0.0	0.0
2.0000	0.0	0.0	0.0	0.0
2.1250	0.0	0.0	0.0	0.0
2.2500	0.0	0.0	0.0	0.0
2.3750	0.0	0.0	0.0	0.0
2.5000	0.0	0.0	0.0	0.0
2.6250	0.0	0.0	0.0	0.0
2.7500	0.0	0.0	0.0	0.0
2.8750	0.0	0.0	0.0	0.0
3.0000	0.0	0.0	0.0	0.0
3.1250	0.0	0.0	0.0	0.0

DATE: 12/11/75      TIME: 14:11:17      PASS: 45  
 WIND: 10.0      DWD: 2.0      SF6: 10.0      SF6: 10.0  
 T-AVE: 10.0      T-STD: 1.0      T-STD: 1.0      T-STD: 1.0  
 S-F6: 10.0      S-F6: 10.0      S-F6: 10.0      S-F6: 10.0

TIME	10.0	10.0	10.0	10.0
0.0000	0.0	0.0	0.0	0.0
0.1250	0.0	0.0	0.0	0.0
0.2500	0.0	0.0	0.0	0.0
0.3750	0.0	0.0	0.0	0.0
0.5000	0.0	0.0	0.0	0.0
0.6250	0.0	0.0	0.0	0.0
0.7500	0.0	0.0	0.0	0.0
0.8750	0.0	0.0	0.0	0.0
1.0000	0.0	0.0	0.0	0.0
1.1250	0.0	0.0	0.0	0.0
1.2500	0.0	0.0	0.0	0.0
1.3750	0.0	0.0	0.0	0.0
1.5000	0.0	0.0	0.0	0.0
1.6250	0.0	0.0	0.0	0.0
1.7500	0.0	0.0	0.0	0.0
1.8750	0.0	0.0	0.0	0.0
2.0000	0.0	0.0	0.0	0.0
2.1250	0.0	0.0	0.0	0.0
2.2500	0.0	0.0	0.0	0.0
2.3750	0.0	0.0	0.0	0.0
2.5000	0.0	0.0	0.0	0.0
2.6250	0.0	0.0	0.0	0.0
2.7500	0.0	0.0	0.0	0.0
2.8750	0.0	0.0	0.0	0.0
3.0000	0.0	0.0	0.0	0.0
3.1250	0.0	0.0	0.0	0.0

•  
•  
•

INDIVIDUAL LINE KEY:

- line 1: time in PST, pass numbers start at the beginning of each experimental day
- 3: DWD is downwind distance from the source
- 6-: data -each profile has been time center averaged to 18 points
  - SF6 is concentration in PPT
  - all lengths are in meters

#### Step 4 - Multi-modal Gaussian Fits

In plume dispersion modelling, mass distributions are most often described by the familiar Gaussian, or normal, distribution on the horizontal plane. To parameterize these models, then, the measured plume cross-sections must also be approximated in a similar fashion.

The success with which a Gaussian shape approximates the actual measured cross-wind profiles will, of course, vary a great deal. Cross-sections were often skewed, multi-modal, or "square-shaped". This analysis step started by determining the standard deviation of the mass from the mean position. In discrete form,

$$\sigma_Y = \sqrt{\frac{N(C-B^2)}{N-1}} \quad (9)$$

$$C = \frac{1}{T} \sum_{i=1}^N f(x_i)(x_i - x_1)^2 \quad (10)$$

$$B = \frac{1}{T} \sum_{i=1}^N f(x_i)(x_i - x_1) \quad (11)$$

$$T = \sum_{i=1}^N f(x_i) \quad (12)$$

where N is total number of points;

$f(x_i)$  is mass, or concentration, at the  $i$  th point;

$x_i$  is the cross-wind position of the  $i$  th point;

These calculations were performed on both the transformed and untransformed profiles of the previous step. As expected, the transformed width was always smaller than the untransformed value, due to the "peak sharpening" effect of the transfer function.

The next task performed in this analysis was a numerical curve fit to the multi-modal Gaussian model, defined as follows:

$$f(y) = \sum_{i=1}^n P_i \exp \left[ \frac{-(y-\mu_i)^2}{2\sigma_i^2} \right] \quad (13)$$

where  $n$  is the total number of modes;

$\mu_i$  is the cross-wind position of the  $i$ th mode;

$f(y)$  is the model value at position  $x$ ;

$P_i$  is the peak concentration of the  $i$ th mode;

$\sigma_i^2$  is the variance of the  $i$ th mode.

Because no unique solution to this curve fitting exists, the program required interactive decisions for each profile. The user initially decided on the number and cross-wind positions of the modes. The program then selected the concentrations at those positions to be the model's amplitude parameters, and calculated the mode variances necessary to minimize the squared deviations from the fit. The fit and observed profile were then graphically displayed for the user. At this point the user could either accept the fit, or alter his/her initial parameters to achieve a more realistic model. Once satisfied, the user "graded" the profiles subjectively in three categories: skewness, ripple, and overall goodness of fit. Skewness refers to the assymetry of

the individual waveforms associated with each mode. Ripple is the high frequency "noise" introduced to the profiles through the fast Fourier transforms. Goodness of fit judges how well the Gaussian model approximates the observed profile. Table 5 lists the complete set of profiles and grades in each category.

The results from this analysis step are stored in the mass storage data set AIR8. An example of AIR8 is supplied in Figure 8. Examples of observed profiles and their associated analytical forms is shown in Figures 9-12. Figure 9 demonstrates a reasonably well-behaved profile. Figure 10 shows a bimodal distribution. The Figures 10 and 11 data hint at an additional mode in the distribution; however, the programmer decided to ignore the minor peak. Some subjectivity was inevitable in this analysis step. The high frequency components in Figure 12, however, are ripple, produced in the FFT. In this case, the model profile is probably closer to reality than the transformed data.

Table 5

Subjectively Determined Quality Analysis  
of Multi-Modal Gaussian Fits  
to SF6 Cross-Sections

1. Grading System

GRADE							
Test	U	A	B	C	Y	N	
Ripple	--	negligible	amplitude of ripple less than 20% of peak	amplitude of ripple greater than 20% of peak	--	--	
Skewness	Undetermined because of waveform overlap	less than 20% of mass displaced	20-50% of mass displaced	greater than 50% of mass displaced	--	--	
Goodness of fit	--	maximum deviation less than 20%	maximum deviation 20-50%	maximum deviation greater than 50%	--	--	
Aligned	--	--	--	--	model aligned with data mode	model aligned with data mode	model aligned with data mode

2. Subscripts

+ identifies a fourth mode

\*\* this pass from "AIR7" is deleted because maximum concentration is less than 1000 ppt.

DAY 1 (6-21-82)

Pass #	ripple	skewness			goodness of fit			aligned		
		peak 1	peak 2	peak 3	peak 1	peak 2	peak 3	peak 1	peak 2	peak 3
5	A	U	U	U	A	A	C	Y	Y	Y
8	C	C	C	--	C	C	--	Y	--	--
9	C	B	--	--	C	--	--	Y	--	--
11	B	B	--	--	B	--	--	Y	--	--
13	B	B	--	--	B	--	--	Y	--	--
14	C	C	--	--	C	--	--	N	--	--
15	A	B	A	--	B	B	--	Y	Y	--
16	C	C	C	--	C	C	--	Y	N	--
17	C	B	--	--	B	--	--	Y	--	--
18	B	A	--	--	A	--	--	Y	--	--
19	*****									
21	C	C	--	--	C	--	--	Y	--	--
23	C	C	--	--	C	--	--	N	--	--
25	C	C	C	--	A	A	--	Y	Y	--
27	A	B	--	--	A	--	--	Y	--	--
28	B	B	--	--	B	--	--	Y	--	--
29	C	C	C	--	B	C	--	Y	Y	--
30	B	B	--	--	A	--	--	Y	--	--
31	C	U	U	U	A	B	A	Y	Y	Y
32	B	C	--	--	C	--	--	Y	--	--
33	B	U	U	U	A	A	A	Y	Y	Y
34	B	C	--	--	C	--	--	Y	--	--
35	B	U	U	U	C	B	B	Y	Y	Y
36	B	C	--	--	B	--	--	Y	--	--
37	A	C	--	--	B	--	--	N	--	--
38	B	C	--	--	A	--	--	Y	--	--
39	B	U	U	--	C	C	--	Y	Y	--
41	B	C	--	--	B	--	--	Y	--	--
42	B	U	U	U, U+	A	A	A, A+	Y	Y	Y, Y+
43	B	B	C	--	B	C	--	Y	Y	--
44	B	U	U	--	B	B	--	Y	Y	--
45	B	C	--	--	C	--	--	Y	--	--
46	A	A	C	--	C	C	--	Y	Y	--
47	A	A	--	--	A	--	--	Y	--	--
48	A	C	C	--	C	B	-	Y	Y	--
49	B	C	--	--	B	--	--	N	--	--
53	A	C	--	--	C	--	--	N	--	--
54	A	U	U	U	C	A	A	N	N	Y
55	B	C	--	--	B	--	--	N	--	--
56	*****									
57	A	C	--	--	C	--	--	Y	--	--
58	B	U	U	U	C	A	A	Y	Y	N
59	A	C	--	--	B	--	--	Y	--	--
61	B	U	U	U	A	A	C	Y	Y	Y
62	A	C	--	--	B	--	--	N	--	--

DAY 1 (6-21-82)  
(cont'd)

Pass #	ripple	skewness			goodness of fit			aligned		
		peak 1	peak 2	peak 3	peak 1	peak 2	peak 3	peak 1	peak 2	peak 3
63	B	C	--	--	B	--	--	Y	--	--
64	A	C	--	--	B	--	--	M	--	--
65	A	C	--	--	B	--	--	N	--	--
68	A	C	C	--	B	B	--	N	Y	--
69	C	C	--	--	B	--	--	Y	--	--
71	B	B	--	--	B	--	--	N	--	--
72	A	B	--	--	B	--	--	Y	--	--
73	B	C	--	--	A	--	--	N	--	--
75	A	U	U	B	B	B	B	Y	Y	N

DAY 2 (6-22-82)

Pass #	ripple	skewness			goodness of fit			aligned		
		peak 1	peak 2	peak 3	peak 1	peak 2	peak 3	peak 1	peak 2	peak 3
3	B	C	C	--	A	A	--	Y	Y	--
4	C	C	C	--	C	C	--	N	N	--
5	A	B	B	--	C	A	--	N	Y	--
6	A	B	--	--	A	--	--	N	--	--
7	C	C	C	--	A	A	--	Y	Y	--
8	C	C	C	--	C	A	--	Y	Y	--
11	A	C	--	--	A	--	--	Y	--	--
12	A	C	--	--	A	--	--	Y	--	--
13	C	C	C	--	C	C	--	N	N	--
14	B	C	C	--	C	C	--	Y	Y	--
15	A	C	--	--	C	--	--	N	--	--
16	A	B	--	--	A	--	--	Y	--	--
17	C	B	--	--	C	--	--	Y	--	--
18	C	C	--	--	B	--	--	Y	--	--
19	B	C	C	--	B	A	--	Y	Y	--
21	A	C	--	--	A	--	--	Y	--	--
22	B	C	--	--	A	--	--	Y	--	--
23	A	B	--	--	C	--	--	Y	--	--
25	C	C	--	--	C	--	--		--	--
26	C	C	--	--	C	--	--	N	--	--
27	C	U	U	U	A	A	A	Y	Y	Y
29	*****									
30	A	U	U	U	B	A	A	Y	Y	Y
31	C	C	--	--	C	--	--	Y	--	--
32	A	U	U	U	A	A	A	Y	Y	Y
33	C	B	C	--	B	C	--	Y	N	--
34	A	B	--	--	A	--	--	N	--	--
35	B	B	--	--	B	--	--	N	--	--
36	A	C	--	--	A	--	--	Y	--	--
37	B	C	C	--	A	C	--	Y	Y	--
41	B	C	C	C	B	C	B	Y	N	N
42	A	C	--	--	A	--	--	Y	--	--
45	B	C	C	--	B	B	--	Y	Y	--
49	C	C	C	C	B	B	B	Y	Y	Y
51	B	U	U	--	C	B	--	Y	Y	--
52	A	U	U	--	B	A	--	Y	Y	--
53	A	C	C	--	B	A	--	N	N	--
54	A	A	--	--	A	--	--	Y	--	--
55	A	B	B	--	A	A	--	Y	Y	--
56	A	C	B	--	C	A	--	Y	N	--
57	C	U	U	--	A	C	--	N	Y	--
58	A	C	--	--	A	--	--	N	--	--
59	A	U	U	U	A	A	A	Y	Y	Y
61	A	U	U	U,U+	B	B	C,C+	Y	N	N,N+
62	A	U	U	U	A	A	B	Y	Y	N
63	B	C	C	--	A	A	--	Y	Y	--
64	A	U	U	--	B	B	--	Y	Y	--
65	A	U	U	--	A	A	--	Y	Y	--

DAY 3 (6-24-82)

Pass #	ripple	skewness			goodness of fit			aligned		
		peak 1	peak 2	peak 3	peak 1	peak 2	peak 3	peak 1	peak 2	peak 3
1	A	C	--	--	C	--	--	Y	--	--
2	B	B	--	--	A	--	--	N	--	--
3	C	C	C	--	B	B	--	Y	Y	--
6	A	A	U	U	A	A	A	Y	Y	Y
7	A	C	--	--	B	--	--	Y	--	--
8	C	C	--	--	C	--	--	Y	--	--
9	A	B	--	--	B	--	--	N	--	--
10	A	B	--	--	C	--	--	N	--	--
12	A	B	--	--	B	--	--	N	--	--
13	A	U	U	--	C	A	--	Y	N	--
14	A	C	--	--	C	--	--	N	--	--
15	B	C	--	--	A	--	--	N	--	--
16	A	A	--	--	A	--	--	Y	--	--
17	A	A	A	--	A	A	--	Y	Y	--
18	B	A	C	--	A	A	--	Y	N	--
19	A	B	--	--	C	--	--	N	--	--
21	B	B	--	--	B	--	--	Y	--	--
22	A	U	U	--	A	B	--	Y	N	--
23	A	B	--	--	C	--	--	Y	--	--
24	*****									
25	A	U	U	--	A	A	--	N	Y	--
26	C	A	--	--	C	--	--	Y	--	--
28	A	B	--	--	A	--	--	Y	--	--
29	B	B	--	--	B	--	--	Y	--	--
30	A	C	--	--	C	--	--	Y	--	--
31	B	C	--	--	C	--	--	N	--	--
32	A	C	C	--	B	B	--	Y	Y	--
33	C	A	--	--	C	--	--	Y	--	--
36	A	B	B	--	A	A	--	Y	N	--
37	A	B	--	--	C	--	--	Y	--	--
38	A	C	C	--	B	A	--	Y	Y	--
39	A	A	C	--	A	A	--	Y	Y	--
41	A	C	--	--	A	--	--	Y	--	--
42	*****									
43	A	C	A	--	A	A	--	Y	Y	--
45	A	B	--	--	C	--	--	Y	--	--
46	B	A	C	--	A	B	--	Y	Y	--
47	A	B	B	--	B	A	--	Y	Y	--
48	A	A	A	--	A	A	--	Y	Y	--
49	A	C	--	--	C	--	--	Y	--	--
53	A	?	?	--	A	A	--	Y	Y	--
54	A	A	C	--	A	A	--	Y	Y	--
55	A	C	--	--	A	--	--	Y	--	--
56	A	B	B	--	A	A	--	Y	Y	--
57	A	B	A	--	A	A	--	Y	Y	--
59	A	C	B	--	C	B	--	Y	Y	--
61	C	C	C	--	B	C	--	N	Y	--

DAY 3 (6-24-82)  
(cont'd)

Pass #	ripple	skewness			goodness of fit			aligned		
		peak 1	peak 2	peak 3	peak 1	peak 2	peak 3	peak 1	peak 2	peak 3
62	B	C	--	--	B	--	--	Y	--	--
63	A	C	U	--	A	C	--	Y	Y	--
64	A	A	--	--	C	--	--	Y	--	--
66	A	C	--	--	C	--	--	Y	--	--
70	A	U	U	--	A	A	--	N	Y	--
72	A	C	B	--	C	B	--	Y	Y	--
73	A	C	C	--	A	A	--	Y	Y	--
74	A	A	--	--	A	--	--	Y	--	--
75	A	B	C	--	A	C	--	Y	Y	--
76	B	B	--	--	B	--	--	N	--	--
77	C	A	--	--	C	--	--	Y	--	--
80	A	A	--	--	B	-	--	Y	--	--
81	B	C	--	--	C	--	--	Y	--	--

DAY 4 (6-25-82)

Pass #	ripple	skewness			goodness of fit			aligned		
		peak 1	peak 2	peak 3	peak 1	peak 2	peak 3	peak 1	peak 2	peak 3
1	A	B	--	--	A	--	--	Y	--	--
2	B	B	--	--	A	--	--	Y	--	--
3	A	B	--	--	A	--	--	Y	--	--
4	C	C	--	--	C	--	--	Y	--	--
6	A	A	--	--	A	--	--	Y		--
7	A	B	--	--	A	--	--	N	--	--
8	A	C	--	--	A	--	--	Y	--	--
9	A	B	--	--	A	--	--	Y	--	--
10	A	C	--	--	A	--	--	N	--	--
11	A	B	--	--	C	--	--	Y	--	--
12	B	C	--	--	C	--	--	Y	--	--
13	A	U	U	--	A	A	--	Y	Y	--
14	A	C	--	--	C	--	--	N	--	--
15	A	C	--	--	C	--	--	N	--	--
16	*****									
17	A	C	C	--	C	C	--	Y	--	--
19	A	B	--	--	A	--	--	N	--	--
21	B	C	--	--	C	--	--	Y	--	--
22	A	C	--	--	B	--	--	N	--	--
23	A	C	--	--	C	--	--	Y	--	--
24	A	B	--	--	A	--	--	N	--	--
25	C	C	--	--	C	--	--	N	--	--
26	C	C	--	--	A	--	--	N	--	--
27	B	C	C	--	A	A	--	N	N	--
28	A	U	U	U	A	A	A	Y	Y	Y
29	A	U	U	--	A	A	--	Y	Y	--
30	A	U	U	U	A	A	A	Y	Y	Y
31	A	C	C	--	C	A	--	Y	N	--
32	A	C	--	--	A	--	--	N	--	--
33	*****									
34	B	U	U	U	A	A	A	Y	Y	Y
35	A		--	--	C	--	--	Y	--	--
36	A	B	--	--	B	--	--	N	--	--
37	A	B	--	--	A	--	--	Y	--	--
38	A	A	--	--	A	--	--	Y	--	--
39	A	U	U	--	A	A	--	Y	N	--
41	A	U	U	--	C	A	--	Y	Y	--
42	A	U	U	--	A	A	--	Y	N	--
43	A	B	--	--	A	--	--	N	--	--
44	B	B	--	--	C	--	--	Y	--	--
45	B		B	--	--	C	--	--	Y	--
46	A	C	--	--	A	--	--	Y	--	--
48	A	U	U	--	A	A	--	Y	N	--
49	A	U	U	--	A	A	--	Y	Y	--
50	A	U	U	--	B	B	--	Y	Y	--
51	A	A	--	--	A	--	--	Y	--	--
52	A	U	U	--	A	A	--	Y	Y	--

DAY 4 (6-25-82)  
(cont'd)

Pass #	ripple	skewness			goodness of fit			aligned		
		peak 1	peak 2	peak 3	peak 1	peak 2	peak 3	peak 1	peak 2	peak 3
53	A	C	C	--	C	A	--	Y	Y	--
54	A	U	U	--	A	A	--	Y	Y	--
55	A	B	--	--	A	--	--	N	--	--
57	*****									
58	A	A	--	--	A	--	--	Y	--	--
59	B	B	A	--	B	A	--	Y	Y	--
61	A	C	--	--	A	--	--	Y	--	--
63	A	?	--	--	A	--	--	Y	--	--
64	A	C	C	--	A	A	--	Y	Y	--
65	A	B	--	--	C	--	--	Y	--	--
66	A	C	--	--	A	--	--	N	--	--
67	A	A	A	--	B	B	--	N	Y	--
68	C	C	--	--	C	--	--	N	--	--
69	A	B	--	--	A	--	--	Y	--	--
70	A	C	--	--	C	--	--	N	--	--
73	A	B	--	--	B	--	--	Y	--	--

DAY 5 (6-27-82)

Pass #	ripple	skewness			goodness of fit			aligned		
		peak 1	peak 2	peak 3	peak 1	peak 2	peak 3	peak 1	peak 2	peak 3
1	A	B	--	--	A	--	--	Y	--	--
2	C	C	--	--	B	--	--	N	--	--
4	A	B	--	--	A	--	--	Y	--	--
5	A	B	--	--	A	--	--	N	--	--
7	A	C	--	--	A	--	--	N	--	--
8	*****									
9	A	B	--	--	A	--	--	N	--	--
11	A	B	--	--	A	--	--	Y	--	--
14	A	C	--	--	A	--	--	N	--	--
16	B	B	--	--	B	--	--	Y	--	--
17	A	B	--	--	A	--	--	Y	--	--
18	C	B	--	--	B	--	--	N	--	--
19	A	B	--	--	A	--	--	Y	--	--
22	A	B	--	--	A	--	--	Y	--	--
24	A	B	--	--	A	--	--	N	--	--
25	C	C	C	--	C	C	--	Y	Y	--
26	A	B	--	--	A	--	--	Y	--	--
27	B	C	C	--	A	C	--	Y	N	--
28	A	B	--	--	A	--	--	Y	--	--
29	A	C	--	--	C	--	--	N	--	--
30	A	C	--	--	C	--	--	N	--	--
31	A	A	--	--	A	--	--	Y	--	--
32	A	B	--	--	Y	--	--	N	--	--
33	A	B	--	--	A	--	--	Y	--	--
34	A	U	U	--	A	A	--	Y	Y	--
35	A	U	U	--	A	A	--	N	Y	--
36	A	C	--	--	A	--	--	N	--	--
37	C	C	--	--	C	--	--	Y	--	--
39	A	A	--	--	A	--	--	Y	--	--
41	A	B	--	--	A	--	--	Y	--	--
42	A	B	--	--	A	--	--	Y	--	--
46	A	B	C	--	A	A	--	Y	Y	--
47	A	B	--	--	Y	--	--	N	--	--
48	B	C	--	--	C	--	--	N	--	--
49	A	B	--	--	A	--	--	N	--	--
51	A	B	--	--	A	--	--	Y	--	--
52	A	C	--	--	C	--	--	Y	--	--
53	A	B	--	--	A	--	--	Y	--	--
43	*****									
44	*****									

DAY 6 (6-28-82)

Pass #	ripple	skewness			goodness of fit			aligned		
		peak 1	peak 2	peak 3	peak 1	peak 2	peak 3	peak 1	peak 2	peak 3
1	A	C	--	--	C	--	--	N	--	--
3	B	B	--	--	A	--	--	N	--	--
4	A	C	--	--	B	--	--	N	--	--
6	C	C	C	--	A	A	--	Y	Y	--
9	B	B	--	--	C	--	--	Y	--	--
11	A	C	--	--	A	--	--	N	--	--
13	A	C	--	--	A	--	--	Y	--	--
15	A	C	C	C	A	A	A	Y	Y	Y
17	*****									
18	A	C	--	--	A	--	--	Y	--	--
19	A	C	--	--	A	--	--	N	--	--
21	A	C	--	--	A	--	--	Y	--	--
22	*****									
23	*****									
26	B	B	--	--	C	--	--	Y	--	--
27	A	C	--	--	B	--	--	Y	--	--
29	A	A	A	--	A	A	--	Y	Y	--
30	*****									
33	C	C	--	--	C	--	--	Y	--	--
34	A	A	--	--	A	--	--	Y	--	--
35	A	C	--	--	C	--	--	Y	--	--
36	A	C	--	--	A	--	--	N	--	--
37	B	C	C	--	B	B	--	Y	Y	--
44	A	C	A	--	A	A	--	Y	N	--
46	A	B	--	--	A	--	--	Y	--	--
48	A	C	--	--	A	--	--	N	--	--
53	A	C	--	--	B	--	--	N	--	--
54	A	C	--	--	B	--	--	N	--	--
55	A	B	--	--	A	--	--	Y	--	--
59	A	B	--	--	C	--	--	Y	--	--
60	A	B	--	--	A	--	--	Y	--	--
61	A	B	C	--	A	A	--	Y	Y	--
62	B	C	--	--	C	--	--	Y	--	--
65	B	C	--	--	B	--	--	Y	--	--
66	B	B	--	--	A	--	--	Y	--	--
67	A	C	--	--	A	--	--	Y	--	--
68	A	C	C	--	A	A	--	Y	Y	--
71	A		--	--	B	B	--	N	N	--
72	A	C	--	--	C	--	--	N	--	--
73	A				A	A	A	N	Y	Y
74	A	B	--	--	A	--	--	Y	-	--
77	C	C	C	C, C+	A	A	A, A+	Y	Y	Y, Y+
78	A	C	--	--	A	--	--	Y	--	--

DAY 7 (6-29-82)

Pass #	ripple	skewness			goodness of fit			aligned		
		peak 1	peak 2	peak 3	peak 1	peak 2	peak 3	peak 1	peak 2	peak 3
1	B	C	--	--	C	----	--	Y	--	--
2	A	U	U	--	A	C	--	Y	Y	--
3	B	C	--	--	B	--	--	Y	--	--
4	B	B	--	--	C	--	--	Y	--	--
5	A	B	--	--	C	--	--	Y	--	--
7	B	C	--	--	C	--	--	N	--	--
12	A	C	--	--	C	--	--	Y	--	--
14	B	C	--	--	C	--	--	Y	--	--
18	A	B	--	--	A	--	--	Y	--	--
19	D	C	--	--	C	--	--	Y	--	--
22	*****									
23	A	C	--	--	C	--	--	Y	--	--
24	A	C	--	--	C	--	--	Y	--	--
26	A	U	U	--	A	A	--	N	N	--
27	A	C	--	--	C	--	--	Y	--	--
29	A	C	--	--	A	--	--	Y	--	--
31	A	C	C	A	A	C	A	Y	Y	Y
32	A	B	--	--	A	--	--	N	--	--
33	B	C	--	--	B	--	--	Y	--	--
35	C	C	--	--	C	--	--	Y	--	--
37	A	B	--	--	A	--	--	Y	--	--
39	A	B	--	--	A	--	--	Y	--	--
41	A	C	B	--	A	A	--	Y	N	--
42	A	B	--	--	C	--	--	Y	--	--
43	A	A	--	--	A	--	--	Y	--	--
44	B	B	B	B	B	C	B	Y	Y	Y
45	C	C	--	--	C	--	--	Y	--	--
46	A	B	--	--	A	--	--	Y	--	--
47	B	C	--	--	C	--	--	Y	--	--
48	B	C	--	--	C	--	--	Y	--	--
50	B	C	--	--	C	--	--	Y	--	--
51	B	B	--	--	A	--	--	Y	--	--
52	B	C	--	--	A	--	--	Y	--	--
53	A	C	U	U	C	B	A	Y	Y	N
56	B	C	C	--	A	B	--	Y	Y	--
58	B	C	C	--	A	A	--	Y	N	--
60	A	U	U	U	A	A	A	Y	Y	Y

Figure 8.

EXAMPLE OF ANALYSIS STEP 4 OUTPUT

```
*****  
# 1: DIR is flight heading in degrees  
# 2: DWD is downwind distance from source  
# 3: X0,Y0 are coordinates of plume where "width"  
# 4: total pass mean is in relation to "width" in  
# 5: the direction of the flight heading  
# 6: waveforms refer to the individual modes of the  
# 7: Gaussian fit  
# 8: peak value is in PPT  
*****  
# 1: DIR is flight heading in degrees  
# 2: DWD is downwind distance from source  
# 3: X0,Y0 are coordinates of plume where "width"  
# 4: total pass mean is in relation to "width" in  
# 5: the direction of the flight heading  
# 6: waveforms refer to the individual modes of the  
# 7: Gaussian fit  
# 8: peak value is in PPT  
*****  
# 1: DIR is flight heading in degrees  
# 2: DWD is downwind distance from source  
# 3: X0,Y0 are coordinates of plume where "width"  
# 4: total pass mean is in relation to "width" in  
# 5: the direction of the flight heading  
# 6: waveforms refer to the individual modes of the  
# 7: Gaussian fit  
# 8: peak value is in PPT  
*****
```

INDIVIDUAL LINE KEY:

- line 1: DIR is flight heading in degrees
- 2: DWD is downwind distance from source
- 3: X0,Y0 are coordinates of plume where "width"  
has a value of zero
- 4: total pass mean is in relation to "width" in  
the direction of the flight heading
- 8: waveforms refer to the individual modes of the  
Gaussian fit
- 9: peak value is in PPT

date : 6/25 time : 1217:11 pass 11

43

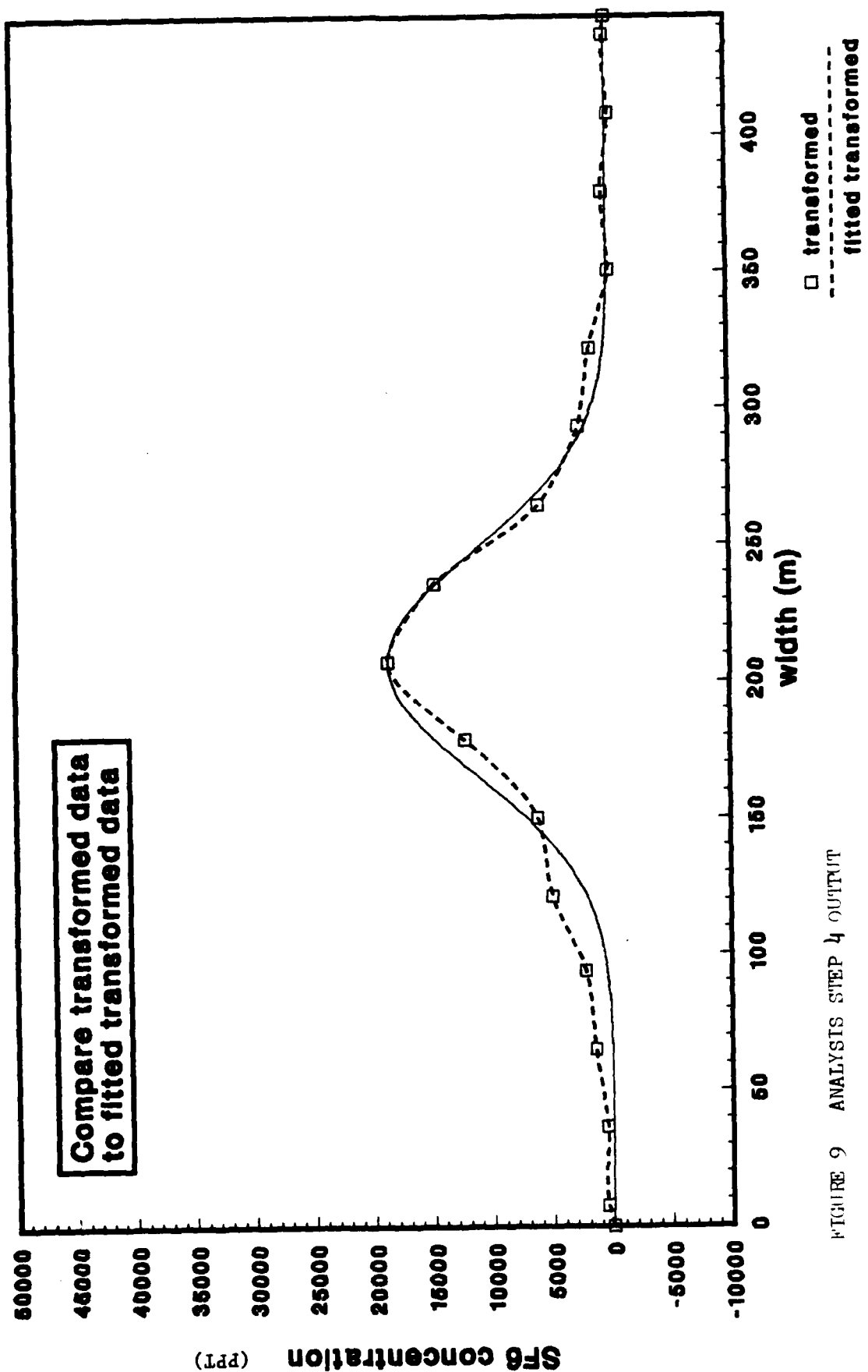


FIGURE 9 ANALYSIS STEP 4 OUTPUT

date : 6/22 time : 1536:22 pass 19

44

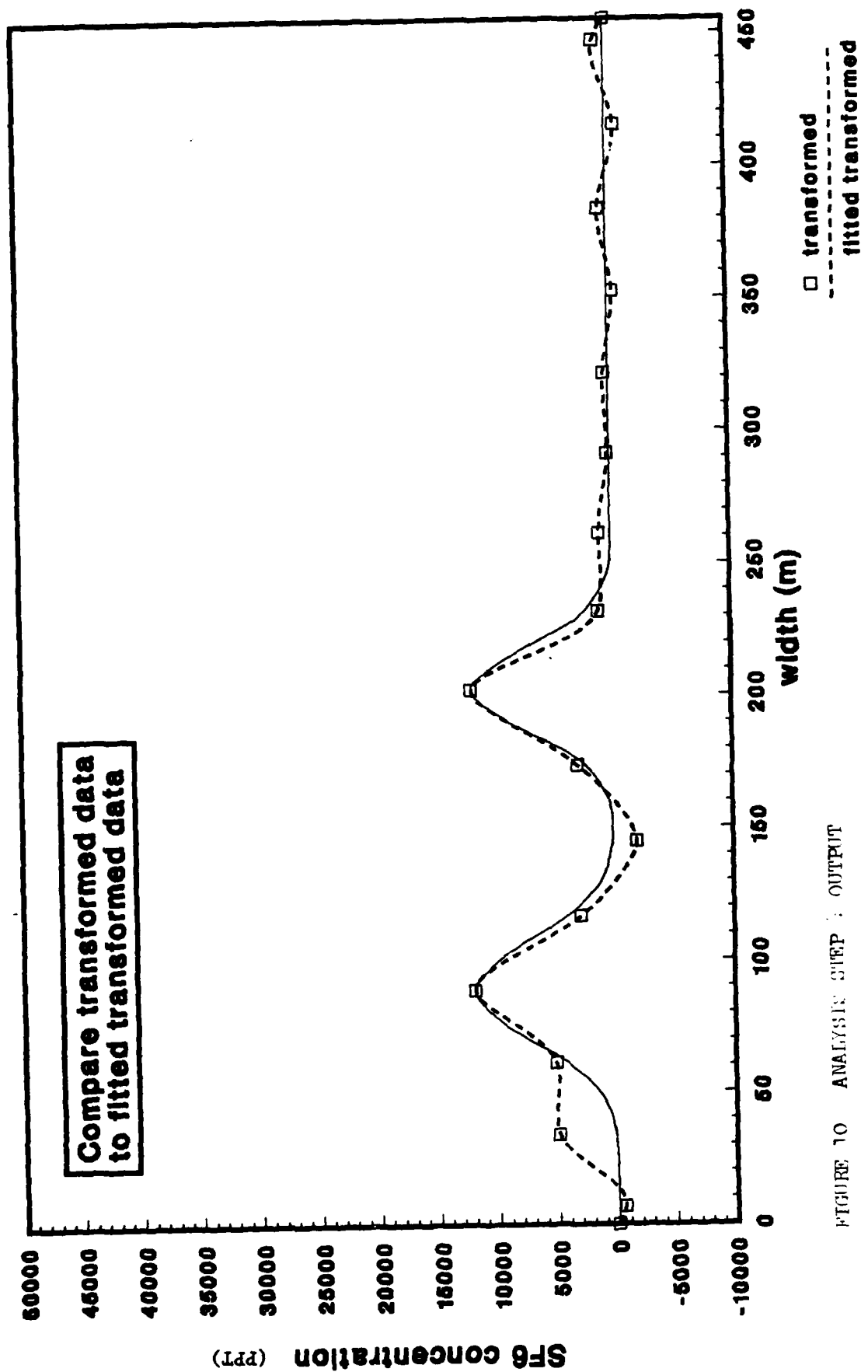


FIGURE 10 ANALYSIS STEP 4 OUTPUT

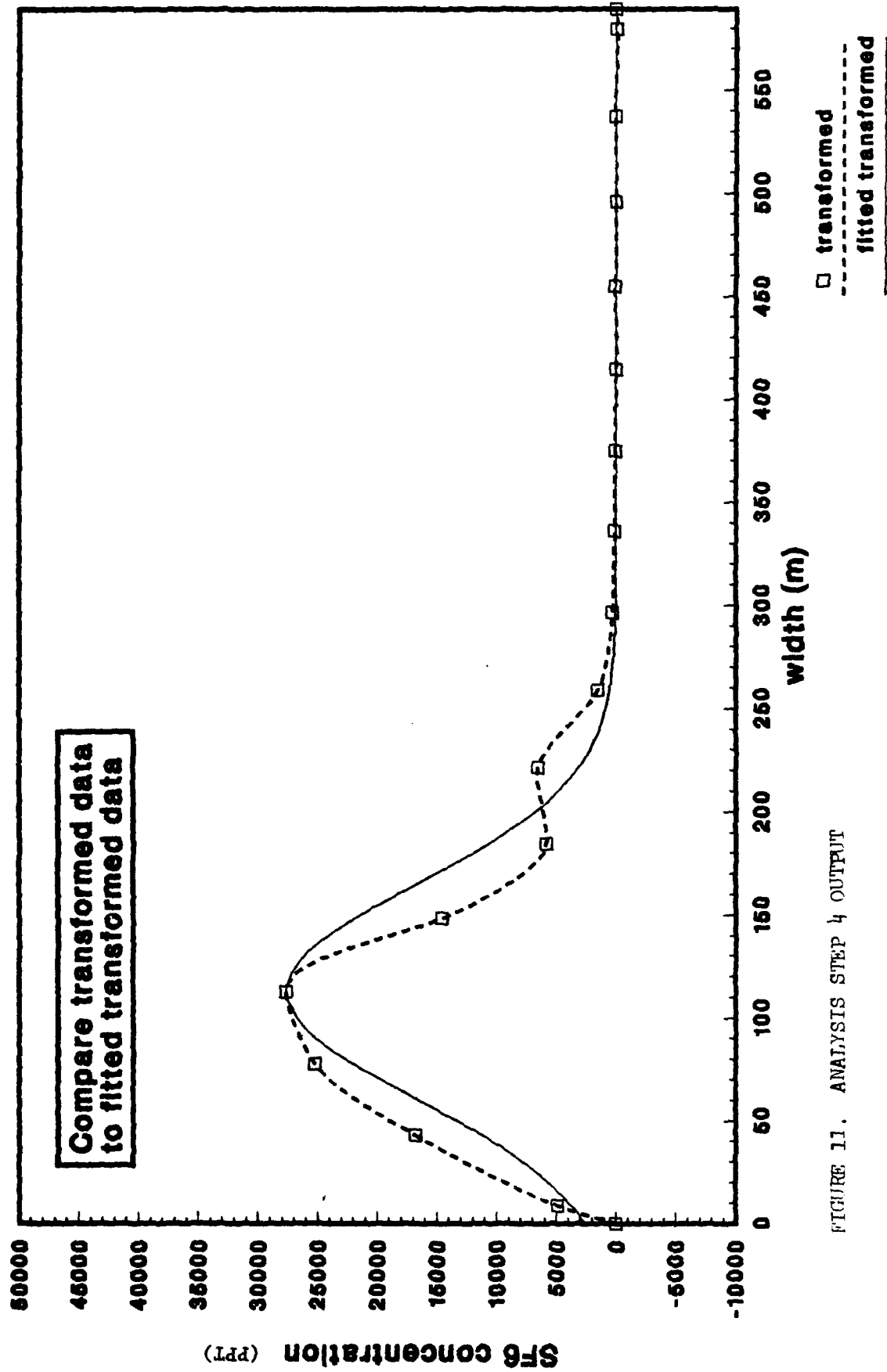


FIGURE 11. ANALYSIS STEP 4 OUTPUT

date : 6/25 time : 1147:30 pass 2

46

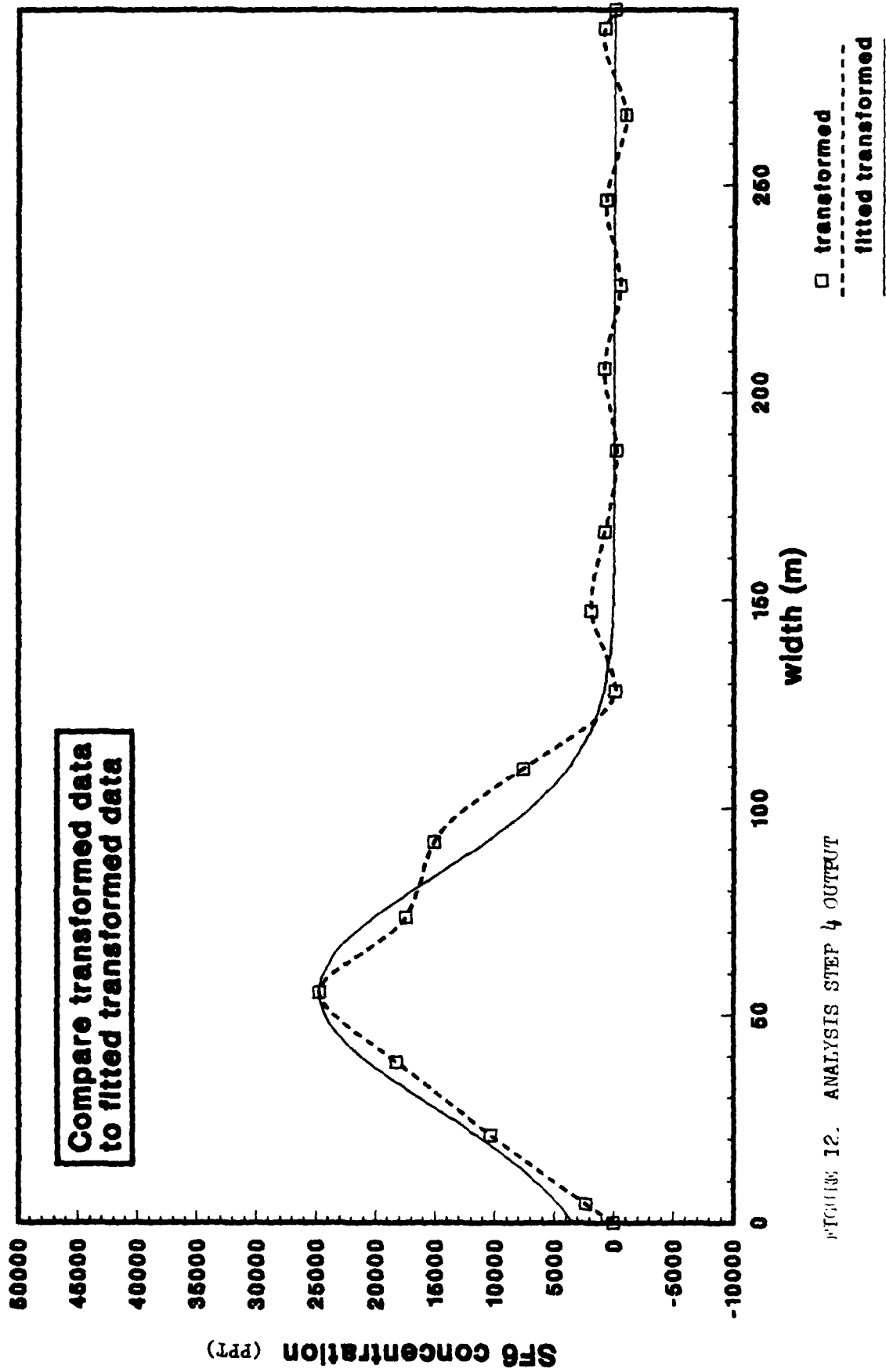


FIGURE 12. ANALYSIS STEP 4 OUTPUT

### Step 5 - Calculation of Hourly Averages

Many dispersion models attempt to predict concentrations expected when averaged over a one-hour period. In order to relate the results of this data set to those of the past, and also in order to satisfy contractual agreements, this analysis step formed one-hour averages of the horizontal and vertical plume growth parameters,  $\sigma_y$  and  $\sigma_z$ . In addition, this step added a header to the data set containing a variety of averaged meteorological quantities.

The basic procedure in this step was to read in half-hour average met data twice, form one-hour average met data, read in tracer data for the current hour, bin the tracer profiles according to range from the release point, and perform the averaging calculations for the plume.

The meteorological data was described in Schacher et al. (1982) and was exclusively collected at the release platform. To account for plume flight-time from the platform, a lag of one-half hour was applied when synchronizing the two-data sets. Even with this adjustment, many problems exist in determining the appropriate meteorology. Due to spatial inhomogeneity, meteorological conditions at the platform become less representative of the average met conditions experienced by the plume as the downwind distance to the aircraft transects increase. Also, meteorological averages tend to differ significantly from hour to hour, implying that stationarity through the one-hour period may be a weak assumption.

For each experimentation day, four range bins were selected, based on the distribution of individual transect downwind distances. An attempt was made to maximize the number of passes in each range bin for all hours, while minimizing the standard deviation of the downwind distances within each range bin. Table 6 lists the range bins for each day.

Table 6  
Range Bins for Hourly Averages of Plume Parameters

	Transect Downwind Distance			
	(m)			
DAY	BIN 1	BIN 2	BIN 3	BIN 4
6-21-82	0-1000	1001-2000	2001-3000	3001-4500
6-22-82	0-1000	1001-2000	2001-3000	3001-4000
6-24-82	0-2000	2001-3000	3001-4000	4001-5000
6-25-82	0-2000	2001-3500	3501-5000	unused
6-27-82	0-2000	2001-3500	3501-5000	5001-6500
6-28-82	0-2500	2501-5000	5001-7500	7501-10,000
6-29-82	0-2500	2501-5000	5001-7500	7501-10,000

One major problem in this analysis step was collecting a sufficiently large number of transects for a range bin during a given hour. Typically, this number was 5 to 12 passes per average.

Discussion of the possibilities and consequences of insufficient sampling will be presented in a later section.

The first averaging calculations performed for each hour-range bin were the average and standard deviation of the bin's downwind distance. The downwind distance (DWD) of a cross-section was interpreted as the straight line distance from the release platform to the plume center. As stated above, the standard deviations of the DWDs for a range bin was minimized to determined range bin boundaries. All DWD standard deviations are less than 200 m.

Five different horizontal plume parameters were calculated for each hour-range bin. Each operated on the ensemble of transects for a bin in a different way. Table 7 gives symbolic definitions used in the discussion that follows.

Table 7  
Definitions for Horizontal Plume Parameters  
Calculated for Each Hour/Range Bin

<u>Symbol</u>	<u>Definition</u>
$\sigma_{yd}$	Mean total standard deviations of the horizontal mass distributions from direct calculations.
$\sigma_{yf}$	Mean total standard deviations from the uni-modal Gaussian fits.
$\sigma_{yw}$	Mean total standard deviations from the uni-modal Gaussian fits weighted by the peak concentration.
$\sigma_{yt}$	Mean total standard standard deviations from the uni-modal Gaussian fits averaged in a fixed cross-wind coordinate system.
$\sigma_{ym}$	Mean standard deviations from the multi-modal Gaussians fits.

$\sigma_{yd}$  was the mean of the standard deviations of the horizontal mass distributions as defined in Equations 9-12, operated on the transformed data. The cross-wind coordinate system was allowed to float in this average. In other words, this average is not affected by plume drift.

$\sigma_{yf}$  is the analytical equivalent of the above. The parameters obtained during the multi-modal Gaussian fits of the previous section were combined to form a single mode fit, and those values averaged. The derivation follows. In continuous form, the mean position of the mass can be defined as the expected value of Y, the "random variable" composed of all y values.

$$E(Y) = \int_{-\infty}^{+\infty} yF(y)dy \quad (14)$$

where Y is the "random variable";

E(Y) is the expected value of Y;

y is the cross-wind position;

F(y) is the density function of y.

The variance is simply the second moment of the distribution, taken about the mean.

$$\sigma_y^2 = E[(Y-\mu)^2] = E(Y^2) - \mu^2 \quad (15)$$

where  $\mu$  is the distribution mean; identically E(Y).

In the case of the multi-modal Gaussian model, the mass, or concentration distribution, is described by Equation 13, repeated:

$$f(y) = \sum_{i=1}^n P_i \exp \left[ \frac{-(y-\mu_i)^2}{2\sigma_i^2} \right] \quad (16)$$

where  $f(y)$  is the concentration at cross-wind position  $y$ ,

$P_i$  is the peak concentration of the  $i$  th mode.

The density function can be formed by simply normalizing Equation 16 by the integrated mass. The mean, or expected value, of  $Y$  is then easily derived as follows:

$$F(y) = \frac{f(y)}{\int_{-\infty}^{+\infty} f(y) dy} = \frac{\sum_{i=1}^n P_i \exp[-(y-\mu_i)^2/2\sigma_i^2]}{\sum_{j=1}^n \sqrt{2\pi} \sigma_j P_j} \quad (17a)$$

$$E(y) = \int_{-\infty}^{+\infty} y F(y) dy = \sum_{j=1}^n \frac{\sigma_j P_j}{\sum_{i=1}^n \sigma_i P_i} \int_{-\infty}^{+\infty} \frac{y \exp[-(y-\mu_j)^2/2\sigma_j^2]}{\sqrt{2\pi} \sigma_j} dy \quad (17b)$$

$$E(y) = \frac{\sum_{i=1}^n \sigma_i P_i \mu_i}{\sum_{i=1}^n \sigma_i P_i} \quad (17c)$$

Rearranging Equation 15 for the  $i$ th mode yields:

$$E_i(Y^2) = \sigma_{yi}^2 + \mu_i^2 \quad (18a)$$

In a similar fashion to the above and using the principle of superposition, it can be shown:

$$E(Y^2) = \frac{\sum_{i=1}^n \sigma_i P_i (\mu_i^2 + \sigma_i^2)}{\sum_{i=1}^n \sigma_i P_i} \quad (18b)$$

Again using Equation 15, the standard deviation of the ensemble profile with n modes is:

$$\sigma_y = \sqrt{\sigma_y^2} = \sqrt{\frac{\sum_{i=1}^n [\sigma_i P_i (\mu_i^2 + \sigma_i^2) - \sigma_i P_i \mu_i^2]}{\sum_{i=1}^n \sigma_i P_i}} \quad (19)$$

$\sigma_{yf}$  was obtained by using Equation 19 for each profile, and averaging all values in each hour-range bin. Results were tested by numerically integrating the same profiles and calculating  $\sigma_y$  as in Equations 9-12. Results were within 1%.

$\sigma_{yw}$  is  $\sigma_{yf}$  weighted by the peak concentration of the member profiles. This parameter is an attempt to bias the mean value toward the cloud width near the plume centroid on the vertical axis, which is ideally at the surface for a surface release. If  $\sigma_y$  is truly independent of height,  $\sigma_{yw}$  should be identical to  $\sigma_{yf}$ .

$\sigma_{yt}$  is defined as the mean total standard deviations from the uni-modal Gaussian fits averaged in a fixed coordinate system.  $\sigma_{yt}$  was obtained in identical fashion to  $\sigma_{yf}$  except each transect was fixed in the cross-wind coordinate system before averaging so that the effects of plume centerline drift are included.

$\sigma_{yt}$  was consistently larger than  $\sigma_{yf}$ . The difference between the values can be interpreted as the degree to which plume meander dominates the hourly averages. In other words, a time-averaged concentration profile can be divided into two components. Plume spread due to relative diffusion, in which there is no fixed axis, is represented by  $\sigma_{yf}$  mean fit. Henceforth, this will be called the diffusive component, and is often referred to as puff diffusion. It is chiefly influenced by turbulence of length scales close to the size of the cloud. Plume spread due to single particle diffusion relative to a fixed axis is theoretically approached in Taylor's (1921) theorem. Plume growth under this theory is influenced by the integrated energy spectrum, or turbulence of all scales.  $\sigma_{yt}$  is representative of this time-averaged quantity. The difference between  $\sigma_{yt}$  and  $\sigma_{yf}$  fit is the time-averaged plume spread due to turbulence of scales either much larger, or much smaller, than the cloud size. The later contributions are negligible. The former turbulence scales tend to move the whole instantaneous plume in a "snake-like" fashion and will hereafter be referred to as the meander component.

The final horizontal plume parameter calculated was  $\sigma_{ym}$ , the mean standard deviations from the multi-modal fits. This quantity was the mean of all the individual mode widths in a floating coordinate system. The origin of multiple modes in an instantaneous profile is yet unexplained; therefore, the significance of this calculation is unknown. This parameter increases only slowly with range, and may, in fact, define the size of coherent turbulent structures.

A parameter calculated during this analysis step closely related to horizontal diffusion was the off-axis position of the mean mass. This is the difference between the actual position of the mean mass and that position calculated from the mean wind vector. The quantity shows any inhomogeneity in the mean wind field, such as a sea breeze's veer with decreasing distance to the shoreline. It also reveals meander produced by motions of time scales longer than the one-hour averaging period.

The vertical standard deviation of the concentration is not measured instantaneously, and therefore must be interpreted from the horizontal cross-sections for each hour-range bin. This was accomplished, when possible, by calculating the cross-wind integrated concentration of each profile, and then performing a single-sided Gaussian fit in the vertical through the data points.

The cross-wind integrated concentration is calculated from the fitted profiles and defined as follows:

$$CWIC_z = \int_{-\infty}^{+\infty} \sum_{i=1}^n P_i \exp[-(y-u_i)^2/2\sigma_{yi}^2] dy \quad (20)$$

$$CWIC_z = \sum_{i=1}^n \sqrt{2\pi} \sigma_{yi} P_i \quad (21)$$

where  $CWIC_z$  is cross-wind integrated concentration  
in ppt-m at a height  $z$ ,

$\sigma_{yi}$  is the standard deviation of the  $i$  th mode in  
a given profile,

$P_i$  is the peak concentration of the  $i$  th mode,

$\mu_i$  is the mean position of the  $i$  th mode.

The model from which  $\sigma_z$  was estimated is:

$$CWIC_z = CWIC_0 \exp \left[ -\frac{z^2}{2\sigma_z^2} \right] \quad (22)$$

where  $\sigma_z$  is the vertical standard deviation of mass.

By linear regression of  $\ln(CWIC_z)$  versus  $z^2$ ,  $\sigma_z$  becomes a  
function of only the slope, while  $CWIC_0$  is a function of the  
intercept as follows.

$$\sigma_z = \sqrt{(2a)^{-1}} \quad (23)$$

$$CWIC_0 = \exp(b) \quad (24)$$

where  $a$  is the slope of the  $\ln(CWIC_z)$  vs  $z^2$  line;

$b$  is the intercept of the line.

Errors in the proposed model presented in Equation 22 can be  
introduced by either a differing vertical shape of the concentra-  
tion profile or a non-negligible deposition of  $SF_6$  onto the sea  
surface. The profile shape was examined by visual inspection of  
the  $\ln(CWIC_z)$  vs  $z^2$  plots. The scatter of the points about the  
regression line appeared to be unbiased in the vertical for the  
cases examined, indicating that the  $\exp(-z^2)$  model was reasonable.

The possibility of mass loss was examined by comparing the ground-level cross-wind integrated concentration predicted by Equation 24 to the value forced by the source emission rate. The Gaussian plume model requires:

$$CWIC_{G*} = \sqrt{\frac{2}{\pi}} \left( \frac{Q}{\sigma_z u} \right) \quad (25)$$

where  $CWIC_{G*}$  is ground-level cross-wind integrated concentration predicted by the Gaussian plume model

$Q$  is the emission rate, 25 lb. SF6/hr,

$\sigma_z$  is the range-dependent vertical plume parameter,

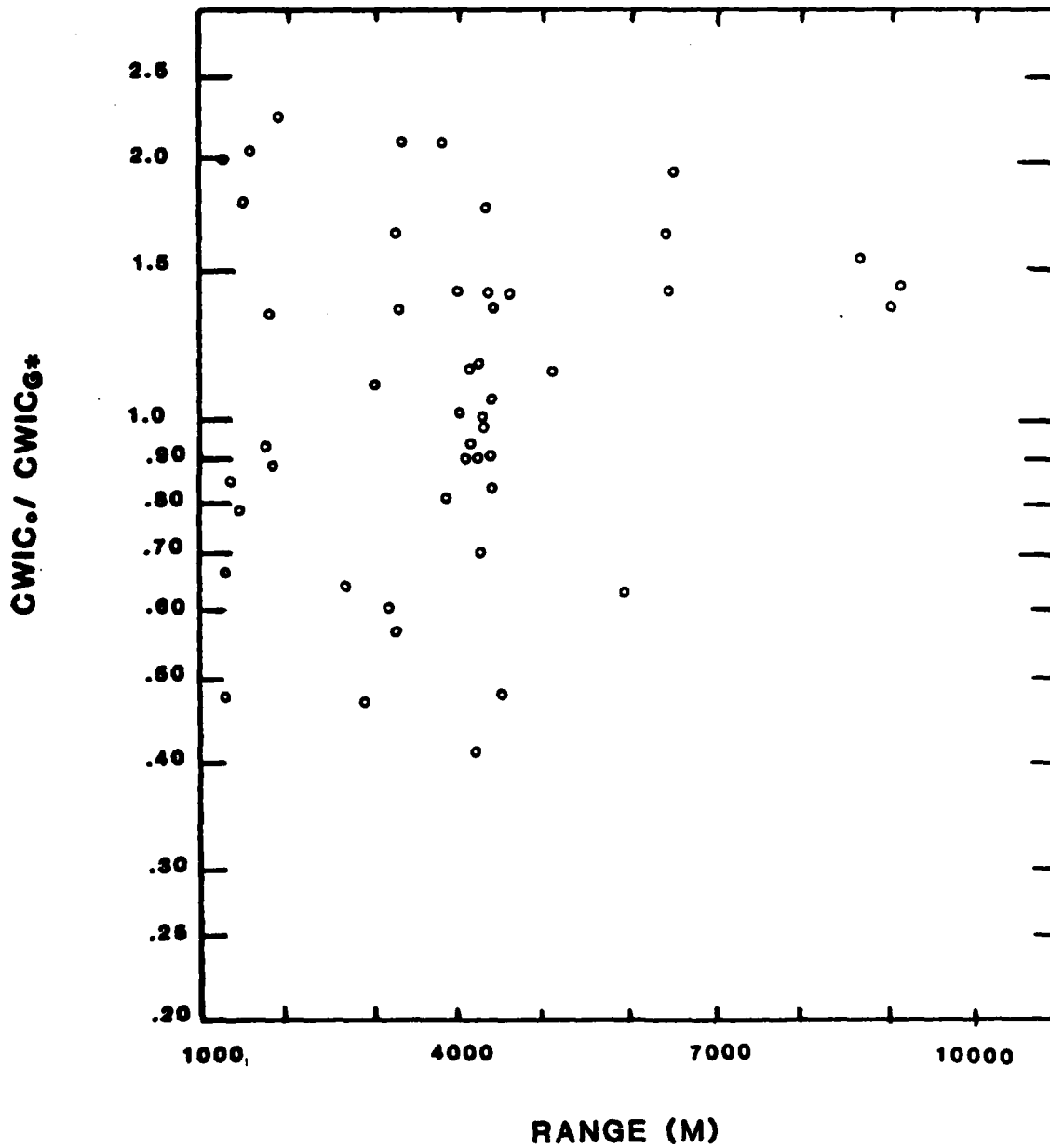
$u$  is the mean wind speed.

Figure 13 shows the ratio of the two values of ground-level cross-wind integrated concentration as a function of range. Ideally, this ratio should be 1 for mass balance. Most points are within a factor of 2. The points are nicely scattered about the identity ratio, and there appears to be no range dependence from 0-9 km.

Based on these results, this analysis suggests that the hourly averaged  $\sigma_z$  values determined by Equation 22 are reasonable.

**FIGURE 13**

**RATIO OF GROUND-LEVEL SF6 MASS CALCULATED  
BY REGRESSION TO MASS DERIVED  
FROM EMISSION RATE**



**see equation (24-25) for quantity definitions**

Step 6 - Plume Parameters as a Function  
of Range and Stability Class

This analysis step uses the hourly averaged tracer and meteorological data produced in Step 5 to parameterize range-dependent plume growth as a function of commonly obtained shipboard meteorological measurements. This step uses only the fixed fit  $\sigma_y$  in the horizontal plume growth parameterization. Future analysis will concentrate on some of the other forms of the horizontal plume dimension, in order to reduce scatter and examine the effects of averaging time.

This analysis attempts to classify the plume properties on a modification of the well-known Pasquill-Gifford table. (See Gifford [1976]). The original scheme first estimates insolation, based on cloud cover and time of day. Insolation range bin and mean windspeed then determine the appropriate stability class. The scheme essentially makes use of the strong relationship between insolation and buoyancy production of turbulence over land, while relating mean windspeed to mechanical turbulence.

This scheme is not applicable over water because, first of all, buoyancy is only weakly dependent on insolation over the oceans, due to the large specific heat of water. Air-surface temperature differences, the primary factor in buoyant production near the surface, are more often the result of advection of either water or air masses than insolation. Second, while mechanical

mixing is still primarily a function of mean windspeed over the ocean, the analytical form of that relationship is quite different.

In order to find a common link between dispersion over water and land, the fundamental physical mechanisms must be examined. At a given height, dispersion is primarily a function of  $z_0$ , the characteristic surface roughness length; and  $L$ , the Monin-Obukhov length, defined as follows:

$$L = \frac{u_*^3 c_p \rho T}{k g H} \quad (26)$$

where  $u_*$  is the friction velocity,  
 $c_p$  is the specific heat of air at constant pressure,  
 $\rho$  is the air density,  
 $T$  is the absolute air temperature,  
 $k$  is the von Karman's constant,  
 $g$  is the acceleration of gravity,  
 $H$  is the vertical heat flux.

In a now-classic paper by Golder (1972), these quantities have been related to the Pasquill-Gifford stability classes. During the BLM experiments, Schacher et al. (1982b) measured the variables necessary to compute  $z_0$  and  $L$ . Schacher et al. (1982a) developed a modified Pasquill-Gifford classification (referred to as NPS scheme) by relating  $z_0$  and  $L$  to routine meteorological measurements, and examined the behavior of  $\sigma_\theta$ , the standard deviation of the wind direction, as a function of the NPS scheme. The analysis reported here extends this concept one step further,

using the NPS scheme of determining stability class together with actual trace gas measurements to build a family of curves.

The Schacher scheme requires four routine meteorological measurements to define stability class: mean windspeed, relative humidity, air temperature, and sea surface temperature. Three sets of curves, for 50%, 80%, and 95% relative humidity, are used to determine the class. Figure 14 shows the result for 50% humidity. From the air-sea temperature difference and the mean windspeed, an appropriate Pasquill-Gifford stability class is chosen by interpolation between curves. The complete set of algorithms is presented in Table 8. Two important points are, first, under this scheme, stability classes A, F, and G are not represented and second, the scheme breaks down at windspeeds under 2 m/s.

At windspeeds under 2 m/s, unless conditions are highly stable, turbulence, and therefore turbulent diffusion, becomes highly inhomogeneous on a horizontal plane. Defining a stability class in order to define plume spread for a Gaussian dispersion model implies homogeneous, steady-state conditions. Defining stability class A over the ocean is probably unnecessary, and may be inappropriate because it is unlikely the sea surface can supply upward heat flux capable of supporting extreme super-adiabatic conditions. Defining classes F and G, on the other hand, is important for coastal regions. Kristensen et al. (1981) gives many over-water examples where these conditions prevail for extended periods of time. Discussion of this problem is given in Appendix B.

FIGURE 14.

EXAMPLE OF NPS OVER-WATER STABILITY CLASSIFICATION SCHEME

\*\*  $\Delta T$  = air temperature - sea surface temperature

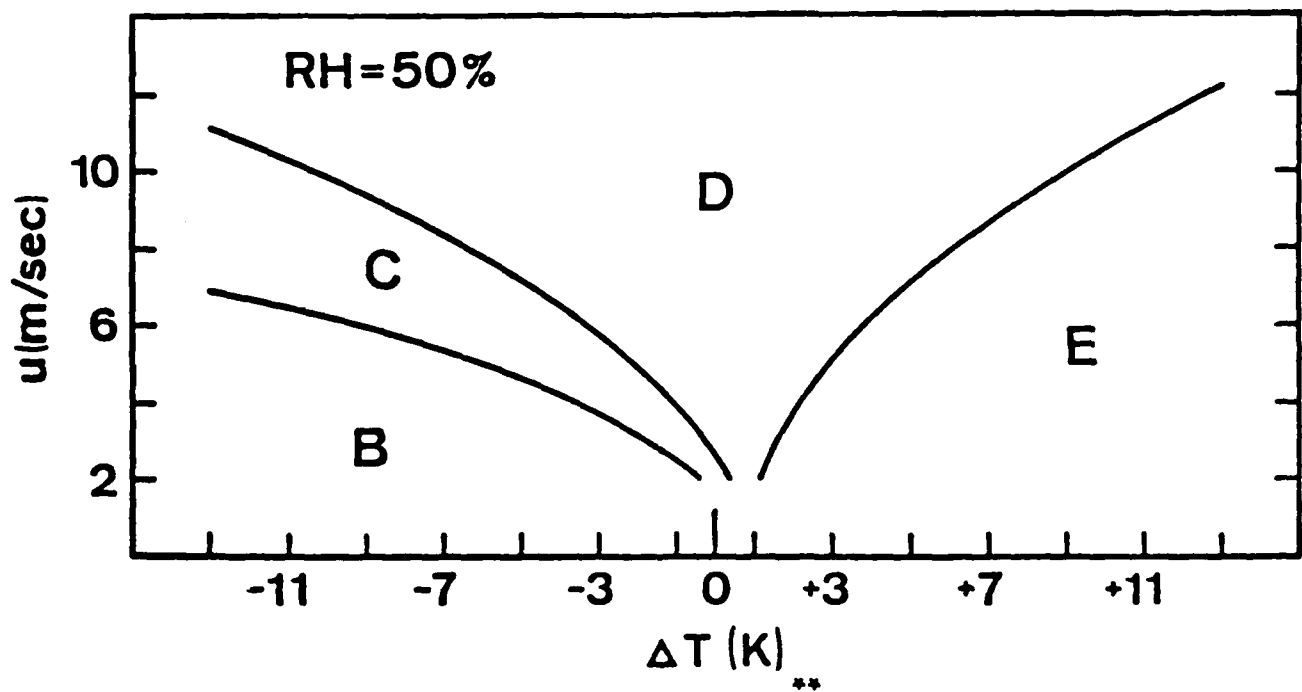


Table 8

P-G Stability Class Scheme  
Adapted to a Marine Boundary Layer

$$U = a_0 + a_1 \Delta T + a_2 \Delta T^2 + a_3 \Delta T^3 + a_4 \Delta T^4$$

where  $a_0, a_1, a_2, a_4$  are constants;

$U$  is windspeed;

$\Delta T$  is (air temperature\* - sea surface temperature) in °C

Relative Humidity	Boundary Line	$a_0$	$a_1$	$a_2$	$a_3$	$a_4$
50%	BC	1.59318	-0.95150	-0.09711	-0.00610	-0.00014
	CD	2.36805	-1.61613	-0.18965	-0.01315	-0.00031
	DE	-0.55452	2.65966	-0.34382	0.02783	-0.00087
80%	BC	1.12799	-1.08521	-0.11388	-0.00707	-0.00016
	CD	1.21695	-2.06787	-0.25450	-0.01708	-0.00040
	DE	0.56149	2.53558	-0.35185	0.03053	-0.00100
95%	BC	1.18368	-0.85413	-0.05274	-0.00248	-0.00005
	CD	1.12545	-1.79684	-0.16237	-0.00869	-0.00017
	DE	0.90463	2.74354	-0.47268	-0.04718	-0.00165

\* Optimum: 10 meter measurement, but any surface layer value is acceptable.

The basic model used to parameterize plume growth for each stability category was

$$\sigma_{y,z}(x) = \sigma_{y,z \text{ ref}} \left( \frac{x}{x_{y,z \text{ ref}}} \right)^{\alpha, \beta} \quad (27a)$$

where  $\sigma_{y,z}(x)$  is the horizontal or vertical standard deviation of the normally distributed mass at range  $x$  ;

$\sigma_{y,z \text{ ref}}$  is a constant for a given stability category representing an appropriate  $\sigma_{y,z}$  at a range  $x_{y,z \text{ ref}}$ ;

$\alpha, \beta$  are constants for a given stability category representing horizontal or vertical plume growth.

For comparison with accepted overland models of similar form,

$x_{y,z \text{ ref}}$  was chosen to be 100 m. To simplify notation, Equation 27a can be expressed as follows:

$$\sigma(x) = bx^c \quad (27b)$$

where  $\sigma(x)$  is either  $\sigma_y$  or  $\sigma_z$  ,

$b$  is either  $\sigma_{y \text{ ref}}/(100)^\alpha$  or  $\sigma_{z \text{ ref}}/(100)^\beta$ ;

$c$  is either  $\alpha$  or  $\beta$ .

The regression analysis was performed in several different fashions (to be discussed) for intercomparison, but all were designed to minimize the mean fractional error, defined as follows:

$$\text{MFE} = \frac{2(P-0)}{P+0} \quad (28)$$

where  $P$  is predicted plume parameter,

$O$  is observed (measured) plume parameter.

Using this error analysis, instead of the usual mean square error, gives logarithmically unbiased results; an over-estimate of  $n \times$  measured value is the same as an underestimate of  $1/n \times$  measured value. This implies that overpredictions are more heavily weighted than underpredictions. This is a desirable trait, since the data set has a lower, but no upper boundary. Also, the standard deviation of the MFE is a measure of the precision (scatter) of the estimate; another useful characteristic. Irwin (1982) gives a similar example of the use of MFE in a sensitivity analysis of overland dispersion models.

Equation 27b contains two unknowns. The coefficient  $b$  essentially represents the initial conditions, or short-range diffusion, which has not been measured directly over the ocean. The exponent,  $c$ , represents the curvature of the scatter plot, or the deviation from linearity of plume growth as a function of range. Regressing  $\ln(\sigma(x))$  versus  $\ln(X)$  and allowing both  $b$  and  $c$  to vary will not yield a unique solution. However, selecting a discrete set of values for either  $b$  or  $c$  will produce a single MFE minimum.

The first regression scheme attempted was to select a discrete set of values for  $c$  and examine the standard deviation of the MFE. In all cases  $\sigma_{MFE}$  varied only slightly, suggesting that there was no preferred combination of  $b$  and  $c$ .

Next,  $c$  was held constant and  $b$  allowed to vary. The value of  $c$  was chosen to be 0.85 for horizontal diffusion. This was based on a sensitivity study of various models as they apply to the Brookhaven over-water oil smoke experiment conducted off Long Island (Michael et al. [1973]). The study suggested the 0.85 value to be appropriate for all stability classes. Over land,  $c$  varies from about 0.80 for Pasquill-Gifford (P-G) class E to 1.00 for classes G-A. In most cases, holding  $c$  constant produced unreasonable values of  $b$ . In other words, the model misrepresented short-range diffusion.

To remedy the problem, the approach was reversed; estimates of short-range diffusion were assumed and the curvature term forced. As previously mentioned, no short-range diffusion data are available. However, statistical theory introduced by Taylor (1921) and applied by Pasquill (1971) and Draxler (1976) allow estimates of short-range diffusion. Specifically, in the horizontal case,

$$\sigma_y(T) = \sigma_v T f_y \left( \frac{T}{t_L} \right) \quad (29)$$

where  $\sigma_v$  is the standard deviation of the cross-wind velocity component;

$T$  is the diffusion time;

$f_y \left( \frac{T}{t_L} \right)$  is a universal function;

$t_L$  is the Lagrangian time scale.

Approximating  $\sigma_v T \approx \sigma_\theta x$

$$\sigma_y(x) = \sigma_\theta x f_y \left( \frac{T}{t_L} \right) \quad (30)$$

where  $\sigma_\theta$  is the standard deviation of the wind direction,  
 $x$  is the downwind distance.

Sheih (1981) has experimentally determined the "universal" function over Lake Michigan for various P-G categories from trajectories of neutrally-buoyant balloons in the surface layer. Sheih (1981) used the model:

$$f_y \left( \frac{T}{t_L} \right) = \left[ 1 + \left( \frac{T}{2t'_L} \right)^{1/2} \right]^{-1} \quad (31)$$

where  $t'_L$  is an "apparent" integral time scale.

Table 9 lists Sheih's experimentally-determined "apparent" integral time scales and Draxler's overland equivalent. Draxler only separated data into stable or unstable; therefore, no "D" value is presented. Notice the large time scale in neutral conditions, representing a large "memory" of an air parcel's trajectory. This is probably a response to synoptic scale disturbances. In non-neutral conditions, the time scale is significantly less than the over-land counterpart.

Equations 30 and 31 can be used with measured values of  $\sigma_\theta$  to obtain horizontal short-range parameters. The  $\sigma_\theta$  values

Table 9.

Sheih's Apparent Integral Time Scales

	P-G CLASS		
	<u>C</u>	<u>D</u>	<u>E</u>
	(all values in seconds)		
horizontal	372 $\pm$ 29(617)	4056 $\pm$ 223	70.3 $\pm$ 3.2(617)
vertical	10.6 $\pm$ 1.1(309)	31.5 $\pm$ 2.1	21.7 $\pm$ 1.3(617)

( ) Draxler's over-land results

Table 10.

Horizontal and Vertical Wind Variance Values\* from  
Central California Air Quality Studies III and IV

P-G Class	# Hrs	1 Hr $\bar{\sigma}_h$	1 Hr $\bar{\sigma}_v$	1 min $\bar{\sigma}_h$	$\sigma_h$ Gifford (76)
B	10	31.0	11.8	7.2	20
C	10	17.3	9.8	7.3	15
D	129	9.1	3.3	2.6	10
E	36	12.6	1.5	2.1	5

\* all values in degrees, and measured at 10 m.

obtained during the 3<sup>rd</sup> and 4<sup>th</sup> Central California experiments used for this procedure are summarized in Table 10. The sample time was one second, and the averaging period was one hour. Also included are the one-minute averaging period values and Gifford's (1976) values for comparison. Note that the over-water values agree with over-land values in all classes except class E. Inspection of the time series and statistical comparison with the well-known "t-distribution" indicates that the large  $\sigma_\theta$  values of class E are statistically significant. These data are probably a large-scale phenomenon, since the one-minute values do not reveal relatively large class E values.

For the vertical case, values of  $\sigma_\phi$ , the standard deviation of the vertical wind direction component, were not measured. They were, however, calculated using surface layer similarity from Binkowski (1978).

$$\sigma_\phi = \frac{u_*}{u} \left[ \frac{\phi_m - z/L}{1.2 f_M} \right]^{1/3} \quad \text{for } z/L > 0 \quad (32)$$

$$\phi_m = 1 + 5 z/L \quad (33)$$

$$f_M = 0.4 [1 + 3.39 z/L - 0.25(z/L)^2] \quad \text{for } 0 < z/L < 2.0 \quad (34)$$

$$f_M = 0.4 [6.78 + 2.39(z/L - 2.0)] \quad \text{for } z/L \geq 2.0 \quad (35)$$

$$\sigma_\phi = \frac{2.89}{u} h^{-.333} \quad \text{for } z/L \leq 0, h > 333 \text{ m} \quad (36)$$

$$\sigma_\phi = \frac{1.14}{u} h^{-.175} \quad \text{for } z/L \leq 0, 25 < h < 333 \text{ m} \quad (36)$$

where L is the Monin-Obukhov length;

u\* is the friction velocity;

u is 10 m windspeed;

h is the inversion height.

As mentioned above, the reference distance used for the short range diffusion parameter,  $\sigma_{y,z \text{ ref}}$ , was 100 m. At this range, Equations 30-31 produce the results presented in Table 11. The minimum and maximum values result from deviations in the "universal function" due to uncertainties in the diffusion time (windspeed) and the apparent integral time scale (error margins in Table 9). Sheih (1981) did not present a value of  $t_L'$  for class B; therefore, values of Table 11 are based on "reasonable"  $t_L'$  values.

Table 11

Calculated  $\sigma_{y,z}$  ref values at 100 m.

Class	min	$\sigma_y$ ref		(m.)	$\sigma_z$ ref	
		mean	max		mean	max
B*	21.65	27.01	32.48	6.17	8.23	10.29
C	24.39	25.90	27.10	6.99	8.70	10.23
D	14.77	15.09	15.41	3.20	3.73	4.19
E	14.35	16.11	17.44	1.34	1.61	1.83

\* only approximate

An interesting aspect of these results is that, for the horizontal case, the class D and E cases are very similar. This is the result of compensating influences of  $\sigma_\theta$  and  $t_L$ ; the smaller  $\sigma_\theta$  values in class D are offset by the larger integral time scale (memory).

With the coefficient term of Equation 28 defined, the exponent can be forced in the regression analysis scheme. The results, the applications, and limitations are presented in the next chapter.

## CHAPTER II - PRELIMINARY RESULTS

### Additional Data Sets

Three additional data sets have been convolved with the data set described in this report (see Table 12). All experiments were conducted with continuous surface releases of the inert gas SF<sub>6</sub>. This implies that the parameterizations derived will be most applicable to a similar release. In addition, Dabberdt et al. (1983) produced some shoreline  $\sigma_y$  and  $\sigma_z$  values from the fourth Central California experiment (BLM IV) which are also incorporated into our data set. The first Gulf of Mexico experiment (GULF I) was conducted during the summer. The warm Gulf water produced the only P-G class B and C conditions that coincided with tracer releases. The third Central California experiment (BLM III) and GULF II were conducted in winter. Cool evening temperatures produced some unstable conditions during BLM III, but these events rarely coincided with tracer releases. GULF II was conducted during a stable, foggy period.

Table 12

#### ADDITIONAL DATA BASES FOR OVERWATER, MEDIUM-RANGE, SURFACE-RELEASE PLUME DISPERSION PARAMETERIZATION

<u>EXPERIMENT</u>	<u>DATE</u>	<u>LOCATION</u>	<u>REFERENCE</u>
Gulf of Mexico Air Quality Study I	Jul 81	Cameron, LA (area)	Dabberdt et al. (1982)
Central California Air Quality Study III	Dec 81	Pismo Beach, CA (area)	Dabberdt, et al. (1983)
Gulf of Mexico II	Feb 82	Cameron, LA (area)	Dabberdt, et al (1982)

The complete set of additional data and method of measurement is supplied in Table 13. Meteorological data is not tabulated, but stability categories were obtained in the manner described in this report.

Table 13

ADDITIONAL 1 HR AVERAGE PLUME PARAMETERS

- all values in meters
- all  $\sigma_z$  values from aircraft transects
- "s" indicates shoreline collectors for  $\sigma_y$
- "a" indicates aircraft transects for  $\sigma_y$
- "b" indicate grab bag samplers from boat

<u>Experiment</u>	<u>Method</u>	<u>Date</u>	<u>HR</u>	<u><math>\bar{\sigma}_y</math></u>	<u><math>\bar{\sigma}_z</math></u>	<u>Range</u>
BLM III	s	12-8-81	13	1225	21.5	6750
	s		14	455	18.5	6880
	s		15	644	15	6700
	s		16	1565	20	7320
	s	12-11-81	13	183	34	6560
	s		14	316	31.5	6630
	s		15	370	24	6660
	s		16	141	27	6660
	s		17	199	-	6820
	s		18	412	-	7190

<u>Exp.</u>	<u>Method</u>	<u>Date</u>	<u>HR</u>	<u><math>\bar{\sigma}_y</math></u>	<u><math>\bar{\sigma}_z</math></u>	<u>Range</u>
BLM III	s	12-17-81	12	--	216.5	6380
	s		13	231	17.5	6510
	s		14	332	-	6380
	s		15	677	116	6630
	s		16	299	39	6860
	s		17	154	22.5	6960
	s		18	387	-	7390
	s	12-14-81	12	194	18	6510
	s		13	200	22.5	6590
	s		14	187	23.5	6530
	s		15	176	12	6600
	s		16	224	-	6740
	s		17	784	-	7310
	s	12-15-81	12	601	79.5	7030
	s		13	346	42	6930
	s		15	723	14.5	6560
	s		16	268	16.5	7010
	s		17	458	35.5	7430
	s		18	812	_____	8290

<u>Exp</u>	<u>Method</u>	<u>Date</u>	<u>HR</u>	<u><math>\bar{\sigma}_y</math></u>	<u><math>\bar{\sigma}_z</math></u>	<u>Range</u>
BLM IV	s	6-21-82	13	559	96	6590
	s		14	148	94.5	6590
	s		15	388	75.5	6640
	s		16	397	76	6670
	s		17	725	—	6590
	s	6-22-82	14	97	11	6280
	s		15	338	41	6380
	s		16	442	42.5	6300
	s		17	241	51.5	6160
	s		18	672	—	6160
	s		19	542	—	6180
	s	6-24-82	12	768	32	6430
	s		13	495	-	6330
	s		14	422	50.5	6280
	s		15	243	48	6250
	s		16	345	-	6290
	s		17	326	-	6590

<u>EXP</u>	<u>Method</u>	<u>Date</u>	<u>HR</u>	<u><math>\bar{\sigma}_y</math></u>	<u><math>\bar{\sigma}_z</math></u>	<u>Range</u>
BLM IV	s	6-25-82	11	117	—	6220
	s		12	219	30	6220
	s		13	260	55	6220
	s		14	239	36.5	6220
	s		15	149	46.5	6240
	s		16	156	—	6260
	s		17	525	—	6430
	s	6-27-82	11	139	—	6820
	s		12	83	—	6610
	s		13	131	—	6670
	s		14	202	34	6630
	s		15	156	39	6650
	s		16	172	32	6720
	s		17	263	32	6640
GULF I	s	7-20-81	13	55	58.5	7019
	b		13	483	—	8661
	s		14	671	—	9275
	b		14	85	—	7480
	s		15	2088	—	8330
	b		15	305	—	6209

<u>EXP</u>	<u>Method</u>	<u>Date</u>	<u>HR</u>	<u><math>\sigma_4</math></u>	<u><math>\sigma_2</math></u>	<u>Range</u>
GULF I	s	7-20-81	16	450	53	8037
	b		16	161	—	5721
	s		17	169	39	9368
	b		17	1492	—	6934
	s	7-23-81	15	870	—	9646
	b		15	354	—	6258
	s		16	498	37	8820
	b		16	203	—	6374
	s	7-27-81	17	233	38.5	8639
	b		17	750	—	5829
	b		19	687	—	6880
	b		19	710	—	5741
	b		20	451	—	7385
	b		20	108	—	6159
	s		21	142	—	7822
	s		21	124	—	5107
	s		13	381	—	8179
	s		13	104	—	5949
	s	7-27-81	14	608	107.5	8055
	s		15	496	115	7872
	b		15	69	—	8501
	s		16	565	—	8058
GULF II	s	2-15-82	13	333	17	4529
	a		13	92	—	2054
	a		13	39	—	1696
	s		14	147	11.5	3992

<u>EXP</u>	<u>Method</u>	<u>Date</u>	<u>HR</u>	<u><math>\bar{\sigma}_y</math></u>	<u><math>\bar{\sigma}_z</math></u>	<u>Range</u>
GULF II	a	2-15-82	14	68	_____	1704
	s		15	679	_____	5170
	s		16	543	_____	5788
	s		17	268	_____	4687
	s		18	125	_____	4507

<u>EXP</u>	<u>Method</u>	<u>Date</u>	<u>HR</u>	<u><math>\bar{\sigma}_y</math></u>	<u><math>\bar{\sigma}_z</math></u>	<u>Range</u>
GULF II	s	2-15-82	19	108	—	4456
	s	2-17-82	13	121	—	6999
	s		14	624	—	6962
	s		15	783	—	7413
	s		16	329	—	7268
	s		17	692	—	6897
	s		18	675	—	7046
	s	2-22-82	12	289	9	7607
	a		12	419	—	4205
	a		12	531	—	4272
	a		12	51	—	4398
	s		13	368	7	7080
	a		13	394	—	3907
	a		13	219	—	3921
	a		13	389	—	4009
	s		14	455	21	6994
	a		14	197	—	3848
	a		14	184	—	3854
	a		14	63	—	3847
	s		15	161	—	7062
	s		16	88	13.5	6957
	a		16	238	—	3846
	a		16	179	—	6401
	a		16	236	—	3847

<u>EXP</u>	<u>Method</u>	<u>Date</u>	<u>HR</u>	<u><math>\bar{\sigma}_y</math></u>	<u><math>\bar{\sigma}_z</math></u>	<u>Range</u>
GULF II	s	2-22-82	17	592	31	6911
	a		17	70	—	3683
	a		17	573	—	4298
	a		17	389	—	3863
	s		18	211	—	7076

<u>EXP</u>	<u>Method</u>	<u>Date</u>	<u>HR</u>	<u><math>\bar{\sigma}_y</math></u>	<u><math>\bar{\sigma}_z</math></u>	<u>Range</u>
GULF II	s	2-23-82	10	498	—	7847
	s		11	238	76.7	7724
	a		11	146	—	4265
	a		11	349	—	4360
	a		11	139	—	4662
	s		12	471	33	8035
	a		12	109	—	4631
	s		12	145	—	4553
	a		12	115	—	4633
	s		13	179	53	7741
	a		13	89	—	4343
	a		13	198	—	4370
	a		13	163	—	4411
	s		14	117	—	7912
	s		15	315	57.7	7984
	a		15	295	—	4545
	a		15	268	—	4546
	a		15	490	—	4044
	s		16	489	46.7	7309
	a		16	395	—	4106
	a		16	313	—	4105
	a		16	490	—	4044
	s		17	107	33.5	7494
	a		17	116	—	7188
	a		17	99	—	7123
	s		18	101	—	7505

<u>EXP</u>	<u>Method</u>	<u>Date</u>	<u>HR</u>	<u><math>\bar{\sigma}_y</math></u>	<u><math>\bar{\sigma}_z</math></u>	<u>Range</u>
GULF II	s	2-24-82	14	186	12	5740
	a		14	163	—	2153
	a		14	139	—	2134
	s		15	186	13	5709
	a		15	123	—	2174
	a		15	105	—	2045
	s		16	83	11.5	6059
	a		16	82	—	2251
	a		16	47	—	2239
	s		17	279	10.5	5822
	a		17	148	—	2160
	a		17	137	—	1975
	s		18	102	—	4722
	s		19	172	—	5155

### Vertical Dispersion Parameters

The encouraging results of the vertical dispersion parameterization are the well-behaved form of  $\sigma_z$  and the distinct difference between classes D and E. The discouraging aspect is that this data contains no class B or C values for  $\sigma_z$ . Figure 15 shows the BLM IV scatter plots and regression curves for classes D and E. Numerical results are presented in Table 14. Also shown is the Turner (1970) overland curves for comparison. The figure shows obvious differences between classes and a general slower overwater growth compared to its overland counterpart. The slower vertical growth is physically realistic when we consider surface roughness. Lower values of  $z_0$  overwater produce smaller vertical velocity fluctuations during stable and neutral conditions, and therefore smaller plume parameters. The additional data sets were not included in the regression analysis. The Gulf data, by the author's admission, showed serious mass balance problems. Both data sets were based on airplane transects over the shoreline, where the internal boundary layer could have altered results. Nonetheless, this data is included in Figure 16 for review, and supports our results. As stated above, tracer data did not coincide with periods of class B or C stability. Meteorological data, however, was logged for 20 complete hours during these conditions (10 hours apiece). Based on the calculated vertical wind variance for these classes, and the well-behaved vertical dispersion in the neutral and stable

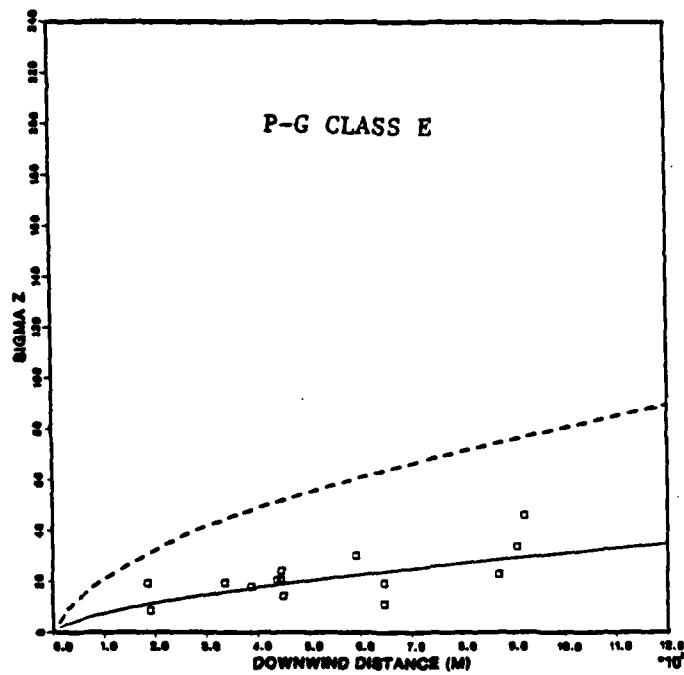
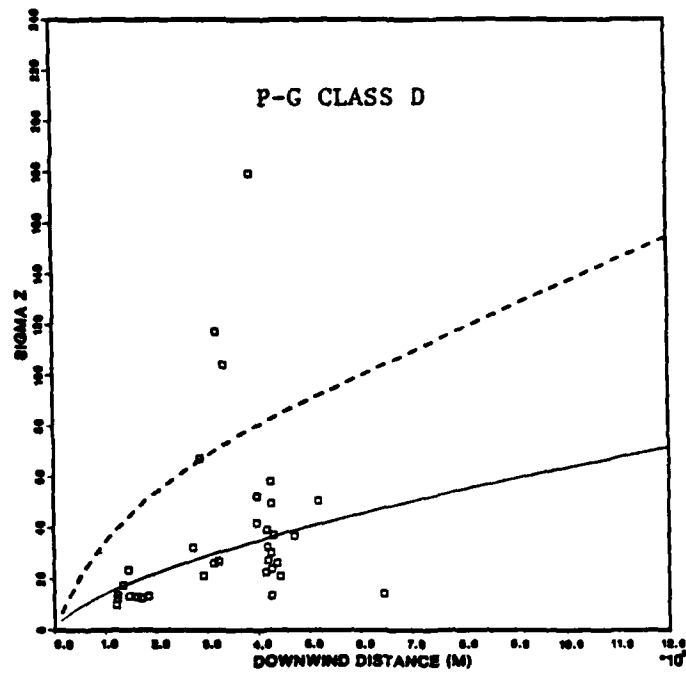


FIGURE 15. ONE HOUR AVERAGE VERTICAL PLUME PARAMETER FROM CCAQ IV

dashed line is Pasquill-Gifford

solid line is table (13)

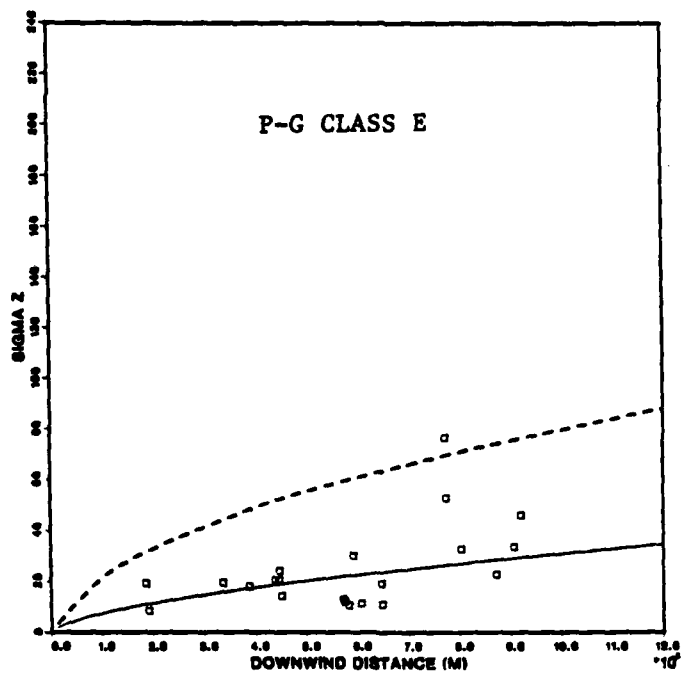
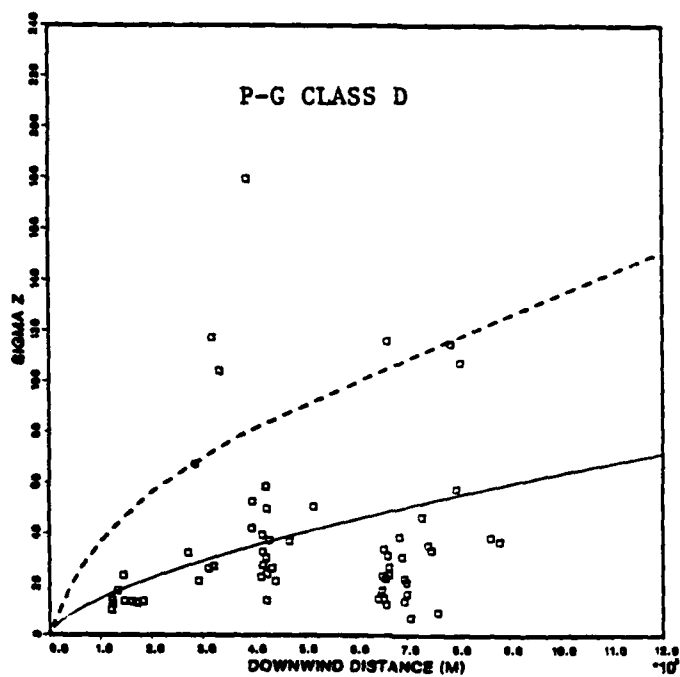


FIGURE 16. ONE HOUR AVERAGE VERTICAL PLUME PARAMETER FROM  
 CCAQ IV AND TABLE (11) DATA SETS  
 dashed line is Pasquill-Gifford  
 solid line is table (13)

Table 14

ONE-HOUR AVERAGE PARAMETERIZATION FOR OVERWATER, SURFACE-RELEASE,  
MODERATE-RANGE\* DISPERSION WITH OVERLAND\*\* COMPARISON

$$\sigma_{y,z}(x) = \sigma_{y,z \text{ ref}} \left( \frac{x}{x_{\text{ref}}} \right)^{\alpha, \beta}$$

$$x_{\text{ref}} = 100 \text{ m.}$$

P-G Category	Over- water $\sigma_{y \text{ ref}}$	Over- land $\sigma_{y \text{ ref}}$	Over- water $\sigma_{z \text{ ref}}$	Over- land $\sigma_{z \text{ ref}}$	Over- water $\alpha$	Over- land $\alpha$	Over- water $\beta$	Over- land $\beta$
B	25.0	19.0	10.0	11.0	0.75a	1.00	0.75a	1.0
C	20.0	12.5	8.0	7.5	0.70a	1.00	0.70a	0.90
D	15.1	8.0	3.2	4.5	0.69	0.90	0.65	0.85
E	16.1	6.0	1.8	3.5	0.65	0.80	0.62	0.80

a insufficient data for verification

\* moderate-range is 0.1-12 km

\*\* Overland values from DTIC (1980)

categories, the shape of the  $\sigma_z$  curve is postulated in Table 14. Verification will proceed as unstable, overwater data become available to the NPS Environmental Physics Group.

#### Horizontal Dispersion Parameters

The hourly averaged horizontal tracer data for P-G classes D and E with regression lines are shown in Figure 17. These results are aesthetically less pleasing than the vertical case because of the increased scatter, but some differences between cases are noteworthy. First, the increased short-range diffusion due to meander for class E, predicted by the theory of the previous section, appears to be realistic when examining the clusters in the 1-2 km range. Second, clusters at greater ranges suggest the overall larger diffusion under class D conditions. The difference is small, however, and the parameterizations of Table 13 reflect this fact. As with the vertical data, P-G classes B and C were insufficiently dense. Ten data points were available in class C, seven in class B, and all data were from GULF I. No regression was attempted on these data, and the values in Table 13 were hypothesized, based on the meteorological ( $\sigma_\theta$ ) data. Verification is needed.

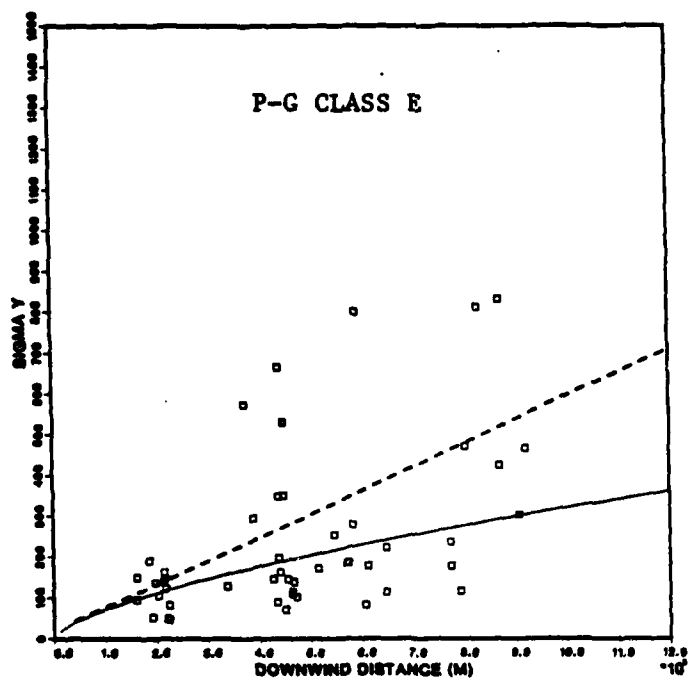
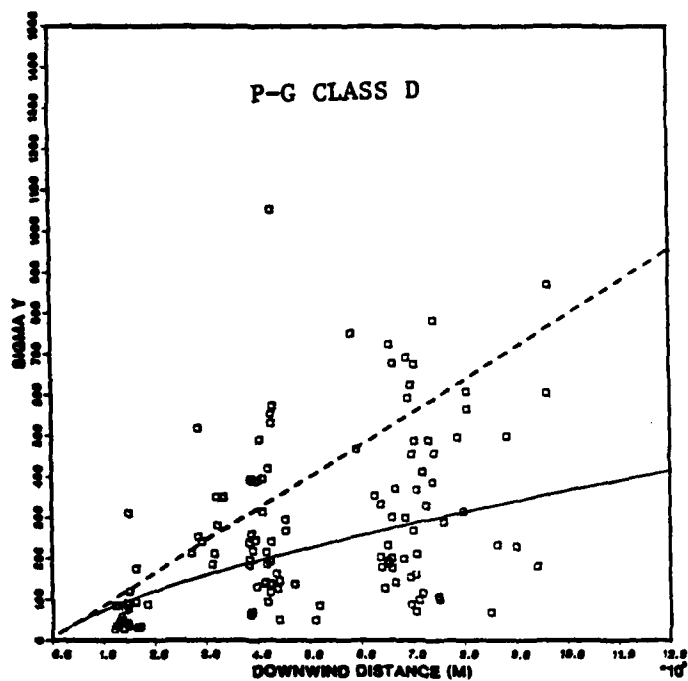


FIGURE 17. ONE HOUR AVERAGE HORIZONTAL PLUME PARAMETERS  
FROM CCAQ IV AND TABLE (11) DATA SETS  
dashed line is Pasquill-Gifford  
solid line is table (13)

As with most tracer data, the points were widely scattered about chosen regression lines. This characteristic feature can be partially attributed to the highly variable nature of turbulence in the atmosphere. Another factor that significantly increases scatter for horizontal data is the large energy in the low frequency part of the horizontal velocity spectra. While a formal spectral analysis of the wind time series was not performed, variance did significantly increase with longer sampling windows, up to one hour. The time series also suggests that this trend would have continued with a larger window. A variety of overland experiments have observed large horizontal wind variance during stable conditions [Hanna (1981), Olesen et al. (1983), Sagendorf and Dickson (1974)]. Spectral analysis by Hanna (1981) indicated a low frequency peak at approximately  $0.5 \text{ hour}^{-1}$ . Olesen et al. (1983) describe large contributions to the energy spectrum at frequencies as low as  $0.35 \text{ hour}^{-1}$ . Kristensen et al. (1981) described increased plume meander in very stable conditions resulting from these low frequency oscillations, and finds an inverse relationship with the mean windspeed (see Appendix B).

Based on the above references, it is not surprising to find a large meander component in the class E  $\sigma_y$  values. It is somewhat unexpected to find a large meander component in near-neutral (class D) stability. These findings are supported in part by Sheih's (1981) large Lagrangian time scales in these conditions, which he has suggested is the result of "large scale motions."

Regardless of the mechanisms involved in the low frequency wind fluctuations, their existence implies that one hour averages are inappropriate for defining horizontal "steady-state" diffusion.

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MEASURED PLUME DISPERSION PARAMETERS OVER WATER VOLUME

2/2

1(U) NAVAL POSTGRADUATE SCHOOL MONTEREY CA

C E SKUPNIEWICZ ET AL. SEP 84 NPS-61-84-812

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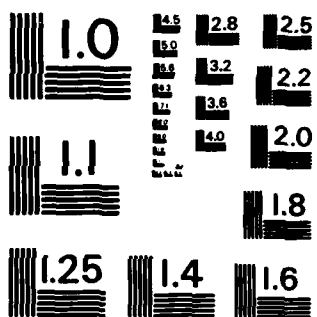
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DTIC



MICROCOPY RESOLUTION TEST CHART  
NATIONAL BUREAU OF STANDARDS-1963-A

## APPENDIX A

### CENTRAL CALIFORNIA AIR QUALITY EXP. IV DATA

The methodology of this analysis was designed in a step-wise fashion to facilitate easy re-analysis. All data sets listed in Table 1A are semi-permanently logged at the NPS Computer Center. Nine-track digital tapes are also available. For the exact data set formats, contact this report's author.

TABLE 1A - TRACER EXPERIMENT DATA SETS AVAILABLE AT NPS

<u>Analysis Step</u>	<u>Program Name</u>	<u>Output Data Set Name</u>	<u>Brief Description</u>	<u>Line#</u>	<u>Ordered content</u>
1	FORMAT	AIR2	raw data	1	code, date time, pass#, data quality index, average altitude
				2	code, elapsed time, #of points, # of transponder 1 pts, # of transponder 2 pts, # bad pts
				3-end	code, mini-ranger#, mini-ranger distance, analyzer output
1	REDUCE	AIR3	calibrated data, rectangular co-ordinates	1	date,time, pass#
					#points
				2	plane heading, air-speed
				3	wind direction, wind speed
				4	standard deviation of output data
				5	cross-wind integrated concentration
				6	transect altitude
				7	distance from release
				8-11	mini-ranger statistics
				12	Release coordinates
				13	
				14-end	elapsed time, running plume width, e-w coordinate, n-s coordinate, concentration
2	XFORM	AIR4	untransformed, arranged data	1	date,time, pass#
				2	altitude, windspeed, wind direction
				3	plane heading, air-speed, distance from release
				4	release coordinates
				5	
				6-24	elapsed time, running plume width, e-w coordinate, n-s coordinate, concentration
2	XFORM	AIR5	transformed, averaged data		Same as AIR4

TABLE 1A - (cont'd)

<u>Analysis Step</u>	<u>Program Name</u>	<u>Output Data Set Name</u>	<u>Brief Description</u>	<u>line#</u>	<u>Ordered Content</u>
3	MINIFIX	AIR6	corrected coordinates-untransformed data		same as AIR4
		AIR7	corrected coordinates-transformed data		same as AIR4
4	FIT	AIR8	multi-modal Gaussian fit	1	null
				2	plane heading, time
				3	altitude, distance from release
				4	e-w coordinate, n-s coordinate
				5	width position of mean mass
				6	standard deviation about mean
				7	total plume width
				8	null
				9	null
				10	peak#1 value, peak#2 value, etc.
				11	peak#1 position, peak#2 position, etc
				12	peak#1 standard dev., peak#2 st. dev., etc.
5	BOTH	AIR9	hourly averages	See Appendix C for complete AIR9 output	
6	BOTH	AIR12	AIR9 condensed	1	date, hour, relative humidity, wind direction, sigma theta
				2	windspeed, air temperature (10m.), sea-surface temperature, 10/L, inversion height
				3	1st average downwind distance (DWD), 1st standard deviation of DWD, 1st # of passes, 1st mean total sigma y
				4	1st mean waveform sigma y

<u>Analysis Step</u>	<u>Program Name</u>	<u>Output Data Set Name</u>	<u>Brief Description</u>	<u>line#</u>	<u>Ordered Content</u>
6	BOTH	AIR7	AIR9 condensed	4 cont.	1st fixed mean total sigma y from fits, 1st off- axis position of mean mass, 2nd average DWD, 2nd st. dev. of DWD
				5	2nd # of passes, 2nd mean total sigma y, 2nd mean waveform sigma y, 2nd fixed mean total sigma y from fits, 2nd off- axis position of mean mass
				6	3rd av. DWD, 3rd standard deviation of DWD, 3rd # of passes, 3rd mean total sigma y, 3rd mean waveform sigma y.
				7	3rd fixed mean total sigma y from fits, 3rd offaxis position of mean mass, 4th average DWD, 4th st. dev. of DWD, 4th # pas.
				8	4th mean total sigma y, 4th mean waveform sigma y, 4th fixed mean total sigma y from fits, offaxis position of mean mass, 1st mean total sigma y from fits
				9	1st weighted mean total sigma y from fits, 1st sigma z, 1st crosswind inte- grated concentration (CWIC), 2nd mean to- tal sigma y from fits 2nd weighted mean total sigma y from fits
				10	2nd sigma z, 2nd CWIC 3rd mean total sigma y from fits, 3rd weighted mean total sigma y from fits,

<u>Analysis Step</u>	<u>Program Name</u>	<u>Output Data Set Name</u>	<u>Brief Description</u>	<u>line#</u>	<u>Ordered Content</u>
6	BOTH	AIR12	AIR 9 condensed	11	3rd CWIC, 4th mean total sigma y from fits, 4th weighted mean total sigma y from fits, 4th sig- ma z, 4th CWIC
				12,13	null

Note: Identically formatted over-water data sets for Central California Air Quality Exp III and the two Gulf of Mexico experiments are also on file.

## APPENDIX B

### OVER-WATER PLUME DISPERSION IN VERY STABLE CONDITIONS

As stated in the main text, very stable conditions are not uncommon over the ocean. These conditions typically occur when the marine boundary layer capping inversions lowers to the sea surface. Under such conditions, the only true measure of stability is the atmospheric temperature lapse rate through the inversion. Dispersion in these conditions departs radically from traditional turbulent diffusion ideas. Kristensen et al. (1981) gives an elaborate theoretical discussion of the physics of dispersion in very stable conditions, identifying the key parameters as averaging time and mean windspeed. Using over-water tracer data at a 20km range, Kristensen found

$$\sigma_y = 3700 T^{1/3} U^{-4/5} \quad (1A)$$

where T is average time;

U is mean windspeed.

This formula is only valid at 20 km, and therefore is of little value to us, but demonstrates the convincingly changed character of diffusion in very stable conditions.

APPENDIX C

COMPLETE HOURLY AVERAGED PLUME PARAMETER INFORMATION FROM THE FOURTH  
CENTRAL CALIFORNIA AIR QUALITY EXPERIMENT

(see Measured Plume Dispersion Parameters Over Water: Volume 2)

— available on request only —

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| 21. | Mr. Peter Guest, Code 63Gp<br>Department of Meteorology<br>Naval Postgraduate School<br>Monterey, California 93943   | 1  |

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| 22. | Mr. Thomas Rappolt<br>Titen Systems Inc.<br>8950 Villa La Jolla Drive<br>Suite 2232<br>La Jolla, California 92038                | 1 |
| 23. | Mr. Dan Goddin<br>ERT<br>975 Business Center Circle<br>Newbury Park, California 21320  | 1 |
| 24. | LCDR James D. Branum<br>2682 Olivestone Way<br>San Jose, California 95132  | 2 |
| 25. | Mr. Lyn Tewscher<br>Titen Systems Inc.<br>8959 Villa La Jolla Drive<br>La Jolla, California 92038                                | 1 |
| 26. | Dr. Welf Aufm Kampe<br>German Military Geophysical Office<br>Mont Royal<br>D-5580 Traben-Trarbach<br>FEDERAL REPUBLIC OF GERMANY | 1 |
| 27. | Office of Naval Research<br>(Attn: CDR S. Colgan, Code 420B)<br>800 N. Quincy Street<br>Arlington, Virginia 22217                | 2 |
| 28. | Office of Naval Research<br>(Attn: Dr. Robert Abbey, Code 425MM)<br>800 N. Quincy Street<br>Arlington, Virginia 22217            | 1 |

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